



Search for the Higgs boson in the $H \rightarrow WW \rightarrow \ell\nu jj$ decay channel at $\sqrt{s} = 7$ TeV with the ATLAS detector

ATLAS Collaboration*

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ABSTRACT

A search for the Standard Model Higgs boson has been performed in the $H \rightarrow WW \rightarrow \ell\nu jj$ channel using 4.7 fb^{-1} of pp collision data recorded at a centre-of-mass energy of $\sqrt{s} = 7$ TeV with the ATLAS detector at the Large Hadron Collider. Higgs boson candidates produced in association with zero, one or two jets are included in the analysis to maximize the acceptance for both gluon fusion and weak boson fusion Higgs boson production processes. No significant excess of events is observed over the expected background and limits on the Higgs boson production cross section are derived for a Higgs boson mass in the range $300 \text{ GeV} < m_H < 600 \text{ GeV}$. The best sensitivity is reached for $m_H = 400 \text{ GeV}$, where the observed (expected) 95% confidence level upper bound on the cross section for $H \rightarrow WW$ produced in association with zero or one jet is 2.2 pb (1.9 pb), corresponding to 1.9 (1.6) times the Standard Model prediction. In the Higgs boson plus two jets channel, which is more sensitive to the weak boson fusion process, the observed (expected) 95% confidence level upper bound on the cross section for $H \rightarrow WW$ production with $m_H = 400 \text{ GeV}$ is 0.7 pb (0.6 pb), corresponding to 7.9 (6.5) times the Standard Model prediction.

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1. Introduction

In the Standard Model (SM), a scalar field with a non-zero vacuum expectation value breaks the electroweak symmetry, gives masses to the W/Z bosons and fermions [1–6], and manifests itself directly as a particle, the Higgs boson [2,3,5]. A primary goal of the Large Hadron Collider (LHC) is to test the SM mechanism of electroweak symmetry breaking by searching for Higgs boson production in high-energy proton–proton collisions. At LHC energies, the Higgs boson is predominantly produced via gluon fusion ($gg \rightarrow H$) and via weak boson fusion ($qq \rightarrow qqH$).

Results of Higgs boson searches in various channels using data up to an integrated luminosity of approximately 5 fb^{-1} have recently been reported by both the ATLAS and CMS Collaborations [7, 8]. The ATLAS analysis excludes a Higgs boson with mass in the ranges 112.9–115.5 GeV, 131–238 GeV and 251–466 GeV while the CMS analysis excludes the range 127–600 GeV at 95% confidence level (CL). Direct searches at LEP and the Tevatron exclude Higgs boson masses $m_H < 114.4 \text{ GeV}$ [9] and $156 \text{ GeV} < m_H < 177 \text{ GeV}$ [10] respectively at 95% CL.

For $m_H \gtrsim 135 \text{ GeV}$, the dominant decay mode of the Higgs boson is $H \rightarrow WW^{(*)}$. For $m_H \gtrsim 200 \text{ GeV}$, the $H \rightarrow WW \rightarrow \ell\nu jj$

channel, where one W boson decays into two quarks leading to a pair of jets ($W \rightarrow jj$) and the other decays into a charged lepton and a neutrino ($W \rightarrow \ell\nu$) where $\ell = e$ or μ , becomes interesting since jets from the Higgs boson decay are, on average, more energetic than the jets from the dominant background ($W + \text{jets}$). An advantage of $H \rightarrow WW \rightarrow \ell\nu jj$ over channels with two final-state neutrinos is the possibility of reconstructing the Higgs boson mass using kinematical constraints to estimate the component of the neutrino momentum along the beam axis.

This Letter describes a search for the SM Higgs boson in the $H \rightarrow WW \rightarrow \ell\nu jj$ channel using the ATLAS detector at the LHC, based on 4.7 fb^{-1} of pp collision data collected at a centre-of-mass energy $\sqrt{s} = 7$ TeV during 2011. The present search supersedes a previous analysis in the same Higgs boson decay channel published by the ATLAS Collaboration [11]. The distribution of the $\ell\nu jj$ invariant mass $m(\ell\nu jj)$, reconstructed using the $\ell\nu$ invariant mass constraint $m(\ell\nu) = m(W)$ and the requirement that two of the jets in the event are consistent with a $W \rightarrow jj$ decay, is used to search for a Higgs boson signal. Feed-down from τ lepton decays is included in this analysis for both background and signal, i.e. $H \rightarrow WW \rightarrow \tau \bar{\nu}_\tau jj \rightarrow \ell \bar{\nu}_\ell \nu_\tau \bar{\nu}_\tau jj$.

The present search is restricted to $m_H > 300 \text{ GeV}$ in order to ensure a smoothly varying non-resonant background. The search is further limited to $m_H < 600 \text{ GeV}$ since, for higher Higgs boson masses, the jets from $W \rightarrow jj$ decay begin to overlap due to the large boost of the W boson, and the natural width of the Higgs

* E-mail address: atlas.publications@cern.ch.

boson exceeds 100 GeV. The best sensitivity to Higgs boson production in this analysis is expected for $m_H \sim 400$ GeV.

2. The ATLAS detector

The ATLAS experiment [12] uses a multipurpose particle detector with forward–backward symmetric cylindrical geometry¹ covering the pseudorapidity range $|\eta| < 2.5$ for charged particles and $|\eta| < 4.9$ for jet measurements. The inner tracking detector (ID) consists of a silicon pixel detector, a silicon microstrip detector, and a transition radiation tracker. The ID is surrounded by a thin superconducting solenoid providing a 2 T axial magnetic field. The superconducting solenoid is surrounded by a high-granularity liquid-argon (LAr) sampling electromagnetic (EM) calorimeter. An iron/scintillator tile calorimeter provides hadronic coverage in the central rapidity range. The end-cap and forward regions are instrumented with LAr calorimeters for both electromagnetic and hadronic measurements. The muon spectrometer surrounds the calorimeters and consists of three large superconducting toroids, each with eight coils, a system of precision tracking chambers, and detectors for triggering.

3. Data and simulation samples

The data were collected using single-muon and single-electron triggers [13]. The single-muon trigger required the transverse momentum (p_T) of the muon with respect to the beam line to exceed 18 GeV; for the single-electron trigger, the threshold varied from 20 GeV to 22 GeV. The trigger object quality requirements were tightened throughout the data-taking period to cope with increasing instantaneous luminosity. For signal electrons satisfying $p_T > 25$ GeV, the trigger efficiency is in the plateau region and ranges between 95% and 97%, depending on the $|\eta|$ of the electron. The muon triggers reaches its efficiency plateau below a signal muon p_T threshold of 20 GeV. The plateau efficiency ranges from about 70% for $|\eta| < 1.05$ to 88% for $1.05 < |\eta| < 2.4$.

Using the ATLAS simulation framework [14], detailed Monte Carlo (MC) studies of signal and backgrounds have been performed. The interaction with the ATLAS detector is modelled with GEANT4 [15] and the events are processed through the same reconstruction chain that is used to perform the reconstruction of data events. The effect of multiple pp interactions in the same and nearby bunch crossings (pile-up) is modelled by superimposing several simulated minimum-bias events on the simulated signal and background events. Simulated MC events are weighted to match the distribution of interactions per beam crossing in the dataset.

4. Object selection

The pp collision vertices in each bunch crossing are reconstructed using the inner tracking system [16]. To remove cosmic-ray and beam-induced backgrounds, events are required to have at least one reconstructed primary vertex with at least three associated tracks with $p_T > 400$ MeV. If multiple collision vertices are reconstructed, the vertex with the largest summed p_T^2 of the associated tracks is selected as the primary vertex.

¹ ATLAS uses a right-handed coordinate system with its origin at the nominal interaction point (IP) in the centre of the detector and the z -axis coinciding with the axis of the beam pipe. The x -axis points from the IP to the centre of the LHC ring, and the y -axis points upward. Cylindrical coordinates (r, ϕ) are used in the transverse plane, ϕ being the azimuthal angle around the beam pipe. The pseudorapidity is defined in terms of the polar angle θ , measured with respect to the z -axis, as $\eta = -\ln[\tan(\theta/2)]$.

Each electron candidate is reconstructed from clustered energy deposits in the EM calorimeter with an associated track. It is further required to satisfy a tight set of identification criteria with an efficiency of approximately 80% for electrons from $W \rightarrow e\nu$ decays with transverse energy $20 \text{ GeV} < E_T < 50 \text{ GeV}$ [17]. While the energy measurement is taken from the EM calorimeter, the pseudorapidity η and azimuthal angle ϕ are taken from the associated track. The cluster is required to be in the range $|\eta| < 2.47$, excluding the transition region between barrel and end-cap calorimeters, $1.37 < |\eta| < 1.52$, and small calorimeter regions affected by temporary operational problems. The track associated with the electron candidate is required to point back to the reconstructed primary vertex with a transverse impact parameter significance $|d_0/\sigma_{d_0}| < 10$ and with an impact parameter along the beam direction of $|z_0| < 1$ mm. Electrons are further required to be isolated: the sum of the transverse energies (excluding the electron itself) in calorimeter cells inside a cone $\Delta R \equiv \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} = 0.3$ around the cluster barycentre must satisfy $\sum(E_T^{\text{calo}})/p_T^e < 0.14$ and the scalar sum of the transverse momenta of all tracks (excluding the electron track itself) with $p_T > 1$ GeV from the primary vertex in the same cone must satisfy $\sum(p_T^{\text{track}})/p_T^e < 0.13$.

Muons are reconstructed by combining tracks in the inner detector and the muon spectrometer. The identification efficiency is measured to be $(92.8 \pm 0.2)\%$ for muons with transverse momentum $p_T > 20$ GeV [18]. Tracks are required to pass basic quality cuts on the number and type of hits in the inner detector. They must lie within the range $|\eta| < 2.4$. The tracks must satisfy the same z_0 cut as electrons and $|d_0/\sigma_{d_0}| < 3$. They must also be isolated, with the sum of the transverse energies (excluding those attributed to the muon itself) in calorimeter cells inside a cone $\Delta R = 0.3$ around the muon satisfying $\sum(E_T^{\text{calo}})/p_T^\mu < 0.14$. Furthermore, the scalar sum of the transverse momenta of all tracks (excluding the muon track itself) with $p_T > 1$ GeV from the primary vertex inside a cone $\Delta R = 0.4$ around the muon must satisfy $\sum(p_T^{\text{track}})/p_T^\mu < 0.15$.

Jets are reconstructed from topological clusters of energy deposited in the calorimeters using the anti- k_t algorithm [19] with radius parameter $R = 0.4$. The reconstructed jet energy is calibrated using p_T - and η -dependent correction factors based on MC simulation and validated with data [20]. The selected jets are required to have $p_T > 25$ GeV and $|\eta| < 4.5$. Jets are considered b -tagged if they satisfy the requirement $|\eta| < 2.8$ and are consistent with having originated from the decay of a b -quark. This latter requirement is determined by a b -tagging algorithm which uses a combination of impact parameter significance and secondary vertex information and exploits the topology of weak decays of b - and c -hadrons. The algorithm is tuned to achieve an 80% b -jet identification efficiency, which results in a tagging rate for light quark jets of approximately 6% [21,22]. The missing transverse momentum and its magnitude E_T^{miss} are reconstructed from calibrated jets, leptons and photons, and take into account soft clustered energy in the calorimeters [23]. Energy deposited by muons is subtracted in the E_T^{miss} calculation to avoid double counting.

5. Event selection

Events are classified based on the number of jets selected in addition to the two jets from the Higgs boson decay candidate. For events to be selected as Higgs boson candidates without an additional jet ($H + 0j$) or with exactly one additional jet ($H + 1j$), the channels which are more sensitive to the gluon fusion process, the following conditions must be met: only one reconstructed lepton candidate (electron or muon) with $p_T > 40$ GeV, no additional leptons with $p_T > 20$ GeV, $E_T^{\text{miss}} > 40$ GeV, and exactly two jets ($\ell\nu jj + 0$ jet sample) or exactly three jets ($\ell\nu jj + 1$ jet sample)

with $p_T > 25$ GeV and $|\eta| < 4.5$. The two jets with invariant mass (m_{jj}) closest to the mass of the W boson are required to satisfy $71 \text{ GeV} < m_{jj} < 91 \text{ GeV}$. One of these two jets must satisfy $p_T > 60$ GeV and the other must satisfy $p_T > 40$ GeV. These two jets are taken as the W boson decay jets and are required to lie within the range $|\eta| < 2.8$, where the jet energy scale is best known (with an uncertainty of 5% or less for $p_T > 40$ GeV, depending on p_T and $|\eta|$ over this range [20]), and have $\Delta R_{jj} < 1.3$ to suppress W + jets background. In order to reduce top quark background, the event is rejected if either of the W boson decay jets is b -tagged.

For the $\ell\nu jj + 2j$ selection ($H + 2j$), which is more sensitive to the weak boson fusion Higgs boson production mode, the following requirements are applied. The charged lepton p_T and the E_T^{miss} must both exceed 30 GeV. There must be at least four jets with $p_T > 25$ GeV and $|\eta| < 4.5$. The two jets with invariant mass closest to the mass of the W boson are required to satisfy $71 \text{ GeV} < m_{jj} < 91 \text{ GeV}$. These jets are labelled as the W boson decay jets. Because of the small signal cross section in this channel, the W boson decay jets are not required to lie within $|\eta| < 2.8$, in order to increase the acceptance. The event is required to satisfy a set of “forward jet tagging” cuts designed to select $qq \rightarrow qqH$ events. The two highest- p_T jets apart from the W boson decay jets are labelled as the “tag” jets, and they are required to be in opposite hemispheres ($\eta_{j1} \cdot \eta_{j2} < 0$). They are also required to be well-separated in pseudorapidity ($\Delta\eta_{jj} = |\eta_{j1} - \eta_{j2}| > 3$). The lepton is required to be between the two tag jets in pseudorapidity. The two tag jets must have large invariant mass ($m_{jj} > 600$ GeV) and there must be no additional jets in the range $|\eta| < 3.2$. The event is rejected if it contains a b -tagged jet.

The $\ell\nu jj + 0/1j$ selection differs from the selection used Ref. [11]. The selection criteria are optimized to improve the expected Higgs boson sensitivity for masses above 300 GeV and require a more complex parameterization of the background shape, as discussed in Section 8.

After the $\ell\nu jj + 0$ and $\ell\nu jj + 1$ selections, the gluon fusion process is expected to contribute approximately 98% and 92% to the total signal yield, respectively, with the remainder primarily due to the weak boson fusion process. After the $\ell\nu jj + 2$ selection, the weak boson fusion process is expected to contribute approximately 68% of the total signal yield, with the remainder primarily due to the gluon fusion process.

6. Expected backgrounds

In both the $\ell\nu jj + 0/1j$ and $\ell\nu jj + 2j$ selections, the background is expected to be dominated by W + jets production. Other important backgrounds are Z + jets, $t\bar{t}$, single top quark, diboson (WW , WZ , ZZ , $W\gamma$ and $Z\gamma$) production, and multijets (MJ) from strong interaction processes that can be selected due either to the presence of leptons from heavy-flavour decays or jets misidentified as leptons.

Although MC predictions are not used to model the background in the Higgs boson search results, a combination of MC and data-driven methods is used to understand the background composition at this intermediate stage. Backgrounds due to W/Z + jets, $t\bar{t}$, and diboson production are modelled using the ALPGEN [24], MC@NLO [25], and HERWIG [26] generators, respectively. Single top production is modelled using AcerMC [27] and single top produced in association with a W boson is modelled with MC@NLO. The small contribution from $W/Z + \gamma$ events is estimated from events simulated using MadGraph/MadEvent [28]. The CT10 parton distribution function (PDF) set [29] is used for the MC@NLO samples, CTEQ6L1 [30] for the ALPGEN and MadGraph samples, and MRSTMCa1 [31] for the AcerMC samples.

The shapes of MJ background distributions are modelled using histograms derived from data samples selected in the same way as for the $H \rightarrow WW \rightarrow \ell\nu jj$ selection, except that the electron identification requirements are loosened and the isolation requirement on muons is inverted. In the loosened selection, electrons satisfying the complete set of identification criteria are not included. Expected contributions from top quark ($t\bar{t}$ and single top) production and electroweak boson (including diboson) production to the MJ shape histograms are subtracted using MC predictions.

To normalize the MJ background contribution in a given channel ($\ell\nu jj + 0j$, $\mu\nu jj + 0j$, $\ell\nu jj + 1j$, $\mu\nu jj + 1j$, $\ell\nu jj + 2j$, $\mu\nu jj + 2j$), a fit to the E_T^{miss} distribution using templates for each background contribution are performed. The E_T^{miss} template is constructed from the loose lepton control sample after the selection is further relaxed by omitting the E_T^{miss} criteria. The normalization of this MJ template and the corresponding template for W/Z + jets taken from MC are fitted to the observed E_T^{miss} distribution in data after the final selection without a E_T^{miss} cut, with other backgrounds estimated using the MC simulation and fixed to their expectation for 4.7 fb^{-1} . The relative contributions from W + jets and Z + jets into the W/Z + jets template are fixed according to the SM cross sections. The scale factors for the MJ and W/Z + jets templates derived from these fits are used to normalize the MJ and W/Z + jets background contributions in comparisons between data and these background expectations.

The MC simulation predicts that W/Z + jets events constitute $(72 \pm 14)\%$ of the total background for $\ell\nu jj + 0/1j$ and $(77 \pm 15)\%$ for $\ell\nu jj + 2j$, while the top quark backgrounds contribute with $(19 \pm 5)\%$ and $(9 \pm 2)\%$ for $\ell\nu jj + 0/1j$ and $\ell\nu jj + 2j$ respectively.

7. WW mass reconstruction

To reconstruct the invariant mass $m(\ell\nu jj)$ of the WW system, the neutrino momentum is required. Its transverse momentum p_T^ν is taken from the measured E_T^{miss} while the neutrino longitudinal momentum p_z^ν is computed using the second degree equation given by the mass constraint $m(\ell\nu) = m(W)$. In the case of two real solutions, the solution with smaller neutrino longitudinal momentum $|p_z^\nu|$ is taken, based on simulation studies. In the case of complex solutions, the event is rejected. This requirement rejects $(20 \pm 1)\%$ of MC signal events at $m_H = 400$ GeV, while for MC W + jets the corresponding rejection is $(30 \pm 1)\%$. These estimates include only statistical uncertainties. Larger fractions of events are rejected in $\ell\nu jj + 1j$ than in $\ell\nu jj + 0j$ independent of lepton flavour. In collision data $(30 \pm 1)\%$ of the events are rejected by this requirement, consistent with the expectations from the W + jets background simulation.

8. Signal and background modelling

The Higgs boson signal is expected to appear as a peak in the $m(\ell\nu jj)$ distribution. Its width, before detector effects, varies from about 10 GeV at $m_H = 300$ GeV to about 70 GeV at $m_H = 550$ GeV. The non-resonant background for the $\ell\nu jj + 0/1j$ channel is modelled by a smooth function of the form $f(x) = [1/(1 + |a(x - m)|^b)] \times \exp[-c(x - 200)]$, where x is $m(\ell\nu jj)$ in GeV and a , b , c , and m are free parameters with the appropriate units. In the $\ell\nu jj + 2j$ channel, the background is modelled by the sum of two exponential functions. The parameters of the fitted function in each of these models are not subjected to any external constraint. The functional form for the background model is well motivated by studies using MC simulation, and is tested by fits to the $m(\ell\nu jj)$ distributions obtained through event selection in the W sidebands, with m_{jj} just below ($45 \text{ GeV} < m_{jj} < 60 \text{ GeV}$) or

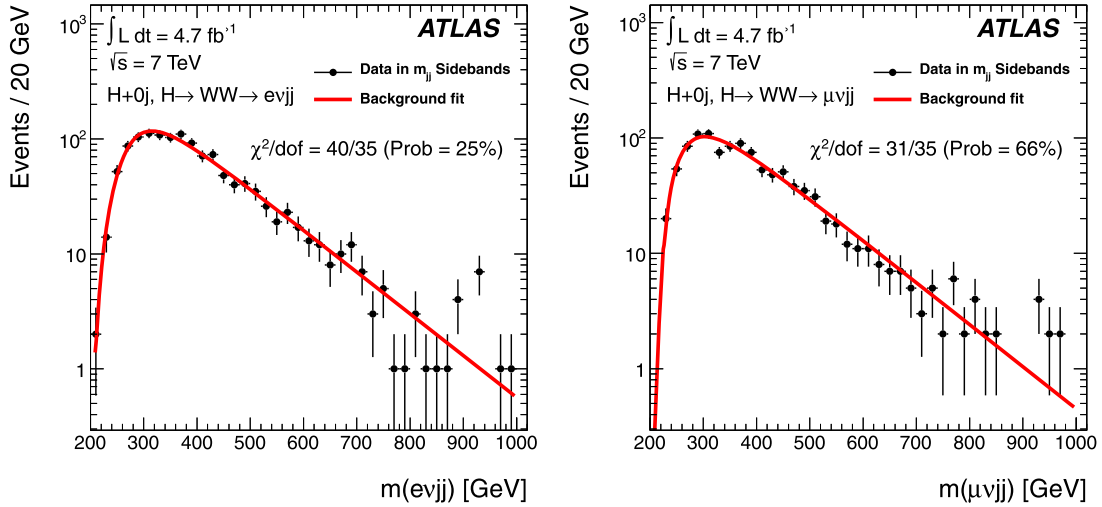


Fig. 1. Fits of the background model described in the text to the reconstructed invariant mass $m(\ell v jj)$ when m_{jj} is in the W sidebands for the $\ell v jj + 0j$ selection. The left (right) figure shows the electron (muon) channel distribution. The χ^2/dof and χ^2 probability of these fits are also shown in the figure.

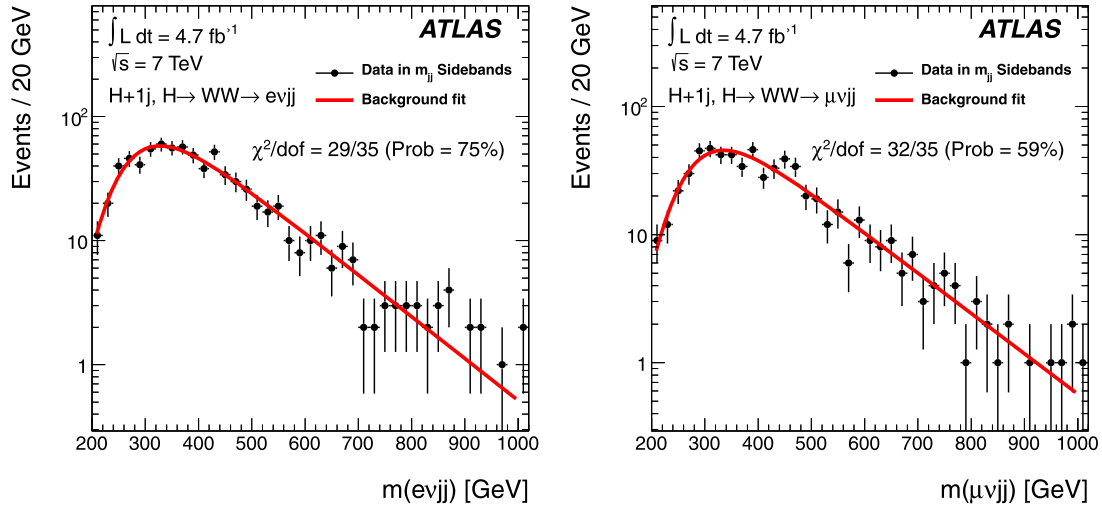


Fig. 2. Fits of the background model described in the text to the reconstructed invariant mass $m(\ell v jj)$ when the m_{jj} is in the W sidebands for the $\ell v jj + 1j$ selection. The left (right) figure shows the electron (muon) channel distributions. The χ^2/dof and χ^2 probability of these fits are also shown in the figure.

just above ($100 \text{ GeV} < m_{jj} < 115 \text{ GeV}$) the W boson peak. Figs. 1 and 2 show fits of the $\ell v jj$ mass to the background model for $\ell v jj + 0j$ and $\ell v jj + 1j$ selections with m_{jj} in the W sidebands. The χ^2 probabilities of these fits are between 25% and 75%, providing support for the background functional form used in this analysis.

MC simulation is used to study the expected Higgs boson contribution to the $m(\ell v jj)$ distributions. Both the gluon fusion and the weak boson fusion signal production processes are simulated using the POWHEG [32,33] event generator interfaced to PYTHIA [34] using MRSTMCAL [31] PDFs and are normalized to the next-to-next-to-leading order cross sections [35] shown in Table 1. The $m(\ell v jj)$ distribution for the expected signal at each hypothesized m_H is modelled using the functional form $1/(a + (x - m_1)^2 + b(x - m_2)^4)$ with parameters (a , b , m_1 and m_2) determined from a fit to the MC simulation of the expected Higgs boson signal. The $m(\ell v jj)$ fractional resolution is $8.8 \pm 1.3\%$ at $m_H = 400 \text{ GeV}$, the uncertainty arising mostly from the E_T^{miss} and jet energy scale as described below, and shows a $1/\sqrt{m_H}$ dependence over the range of this analysis.

Table 1

Cross sections for Standard Model Higgs boson production and the branching ratio (BR) for $H \rightarrow WW \rightarrow \ell v jj$ ($\ell = e$ or μ) as a function of Higgs boson mass m_H . The cross section and its associated uncertainties are described in Ref. [36]. The branching ratio includes $W \rightarrow \tau \rightarrow \ell$, and the uncertainties from the subchannels [37] are added in quadrature with the $H \rightarrow WW$ uncertainty, which is 0.5% below 500 GeV and $0.1m_H^4$ for $m_H \gtrsim 500 \text{ GeV}$.

m_H [GeV]	$\sigma(gg \rightarrow H)$ [pb]	$\sigma(qq \rightarrow H)$ [pb]	$\text{BR}(H \rightarrow \ell^\pm v jj)$
300	2.4 ± 0.4	0.30 ± 0.01	0.237 ± 0.003
400	2.0 ± 0.3	$0.162^{+0.010}_{-0.005}$	0.199 ± 0.002
500	0.85 ± 0.15	$0.095^{+0.007}_{-0.003}$	0.187 ± 0.002
600	0.33 ± 0.06	$0.058^{+0.005}_{-0.002}$	0.191 ± 0.003

9. Systematic uncertainties

The systematic uncertainty due to the background modelling is included by treating the uncertainties on the background model parameters resulting from fits to the data as nuisance parameters in the statistical interpretation of the data. Both the background model and the sum of signal and background models are found to

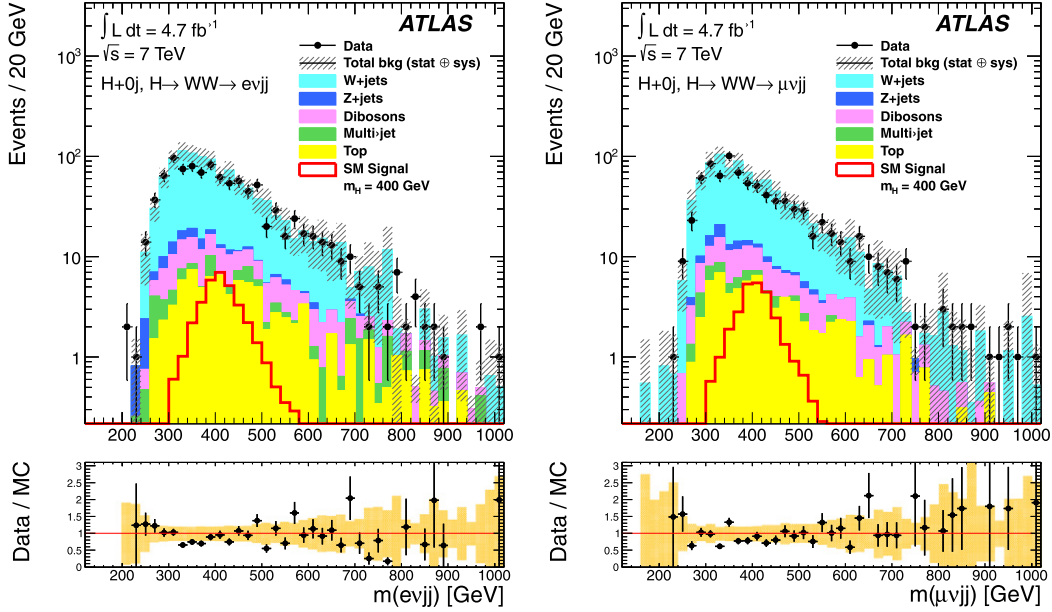


Fig. 3. The reconstructed invariant mass $m(\ell\nu jj)$ in the data and expected backgrounds using MC simulation for the $\ell\nu jj + 0j$ selection. The left (right) figure shows the electron (muon) channel distribution. The expected Higgs boson signal for $m_H = 400$ GeV is also shown. The bottom panels show the data divided by the MC expectation as markers, and the shaded (orange in the web version) region indicates the systematic uncertainty on the background expectation from MC simulation.

be good fits to the data. For $m_H = 400$ GeV, the χ^2 probabilities are 33% and 31% for the background-only and background-plus-signal fits, respectively. Therefore, alternative parameterizations of the background expectation that are consistent with the data will also be consistent with the background model within its uncertainties. This is tested by fitting both the signal region and the sideband regions of the data with two alternative parameterizations that use polynomials of varying order to describe the decreasing background component instead of exponential functions. Differences in the fitted background yield between these parameterizations and the nominal background model are less than 5%, while the uncertainty from the nuisance parameters and statistical uncertainty is 10–12%.

The remaining systematic uncertainties are related to the Higgs boson signal. The fit includes nuisance parameters which account for the uncertainty in the reconstruction efficiency. The trigger efficiencies, the electron and muon reconstruction efficiencies, lepton energy resolution and scale are varied within their uncertainties, giving an uncertainty in the signal efficiency of less than 1%. Varying the jet energy scale [20] within its uncertainties yields an uncertainty of up to 8% in the expected signal in the $\ell\nu jj + 0/1j$ channel for $m_H \geq 400$ GeV. Smearing the jet energies within the uncertainty on their resolutions [38] results in a signal uncertainty of 7% for $m_H = 400$ GeV and 5% for $m_H = 600$ GeV. The reconstructed E_T^{miss} [23] is also affected by the uncertainties on the energy scales and resolutions of reconstructed leptons and jets. The signal uncertainties given above include the propagation of these effects to the reconstructed E_T^{miss} . The propagation to E_T^{miss} adds a small contribution to the overall signal uncertainty. In addition, a 7% uncertainty on the degradation of the E_T^{miss} resolution and scale due to pile-up effects is estimated, which results in a negligible uncertainty on the signal efficiency. The looser selection criteria for the $\ell\nu jj + 2j$ channel result in an 11% uncertainty on the signal efficiency from the jet energy scale at $m_H = 400$ GeV while the uncertainty due to the jet energy resolution is 16%. The uncertainty on the b -tagging efficiency [39] gives a maximum uncertainty of 8% on the signal efficiency and shows no strong dependence on m_H or the selection criteria.

The uncertainties on jet energy resolution and jet energy scale, which also have an impact on E_T^{miss} , lead to systematic uncertainties on the Higgs boson mass resolution (5%) and on the Higgs boson mass scale (2%). These uncertainties are not included since their effect on the fitted Higgs boson yield is considerably smaller than the systematic uncertainty on the signal acceptance due to jet energy scale and resolution.

The Higgs boson signal expectation includes a 3.9% systematic uncertainty due the luminosity determination [40,41] and a 19.4% uncertainty on the predicted Higgs boson cross section [35], taken to be independent of the mass. Off-shell effects and interference between the signal and background processes are discussed in Refs. [35,42,43]. To account for the uncertainties from these effects, an uncertainty of $150\% \times m_H^3$ (m_H in TeV) on the signal cross section is included in the statistical interpretation of the data, where the m_H^3 form is motivated by the scaling of the Higgs boson width with m_H and the normalization factor of 150% is chosen to give $\sim 30\%$ at $m_H = 600$ GeV [35].

10. Results and conclusions

Figs. 3, 4 and 5 show the $m(\ell\nu jj)$ distributions and the ratio of data to background expectation from MC simulation for the six different final states considered in this analysis, along with bands showing the total background uncertainty. The simulated background is not used in the statistical interpretation of the data. Instead, the parameterizations described in Section 8 are used to model the background.

The Higgs boson signal yield in each final state is determined using a binned maximum likelihood fit to the observed $m(\ell\nu jj)$ distribution in the range $200 \text{ GeV} < m(\ell\nu jj) < 2000 \text{ GeV}$. As a check, fits over a smaller range ($200 \text{ GeV} < m(\ell\nu jj) < 1000 \text{ GeV}$) were also performed and the results were found to be consistent with the results presented here.

The difference between data and the fitted background is shown in Fig. 6. The expected signals for $m_H = 400$ GeV and $m_H = 600$ GeV are also shown, each scaled to the 95% CL limit on the production cross section.

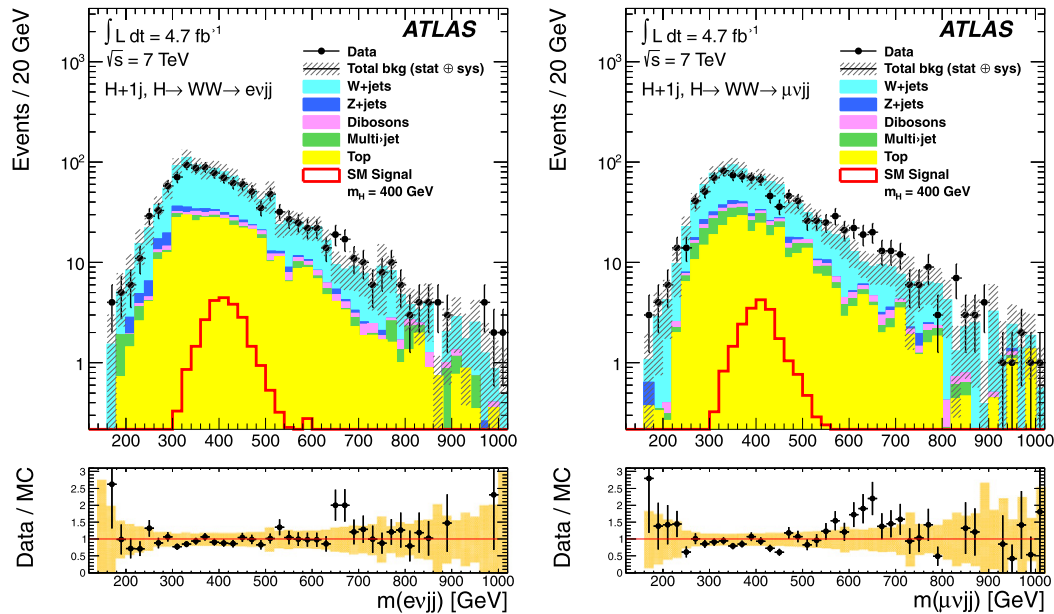


Fig. 4. The reconstructed invariant mass $m(\ell v jj)$ in the data and expected backgrounds using MC simulation for the $\ell v jj + 1j$ selection. The left (right) figure shows the electron (muon) channel distribution. The expected Higgs boson signal for $m_H = 400$ GeV is also shown. The bottom panels show the data divided by the MC expectation as markers, and the shaded (orange in the web version) region indicates the systematic uncertainty on the background expectation from MC simulation.

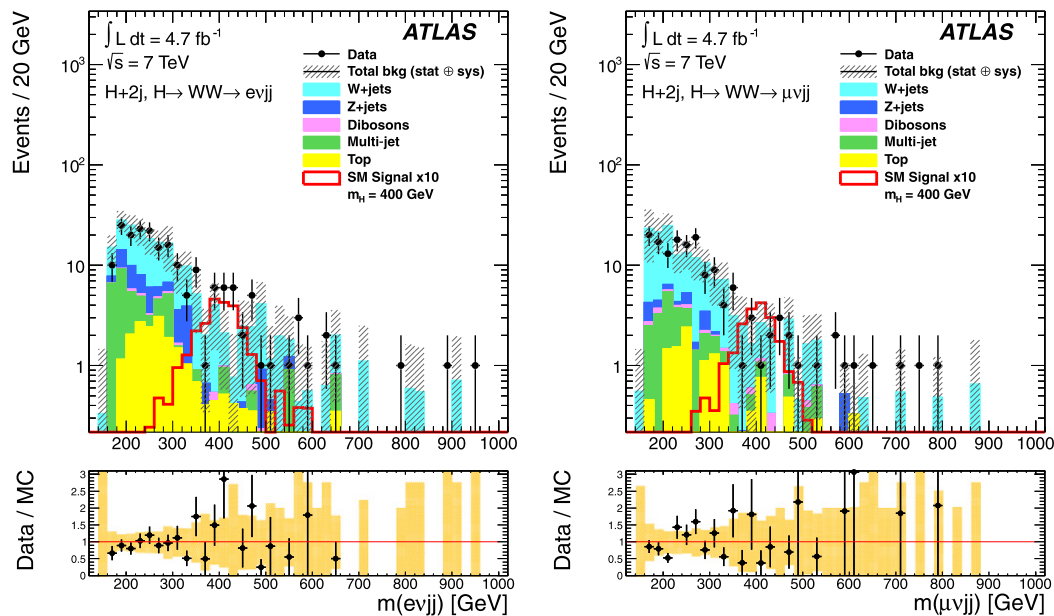


Fig. 5. The reconstructed invariant mass $m(\ell v jj)$ in the data and expected backgrounds using MC simulation for the $\ell v jj + 2j$ selection. The left (right) figure shows the electron (muon) channel distribution. The expected Higgs boson signal for $m_H = 400$ GeV is also shown, scaled up by a factor of 10 for visibility. The bottom panels show the data divided by the MC expectation as markers, and the shaded (orange in the web version) region indicates the systematic uncertainty on the background expectation from MC simulation.

Fig. 6 shows that there is no indication of a significant excess of data above the background model. Limits on SM Higgs boson production are extracted using the profile likelihood ratio [44] as a test statistic and following the CL_s procedure described in Refs. [45,7].

Fig. 7 shows the 95% CL upper bound on the cross section times branching ratio for Higgs boson production with respect to the Standard Model prediction, as a function of m_H . The best sensitivity is reached at $m_H = 400$ GeV, where the 95% confidence level upper bound on the cross section for $H \rightarrow WW$ production using

the combined $H + 0j$ and $H + 1j$ channels is observed (expected) to be 2.2 pb (1.9 pb) corresponding to 1.9 (1.6) times the Standard Model prediction. In the $H + 2j$ channel, which is more sensitive to Higgs boson production via weak boson fusion, the 95% confidence level upper bound on the cross section for $H \rightarrow WW$ production with $m_H = 400$ GeV is observed (expected) to be 0.7 pb (0.6 pb) corresponding to 7.9 (6.5) times the Standard Model prediction. Fig. 8 shows the limits obtained when combining the $H + 2j$ channel with the $H + 0/1j$ channels. Fig. 9 shows the probability p_0 to observe a fluctuation in $300 < m(\ell v jj) < 600$ GeV at least as

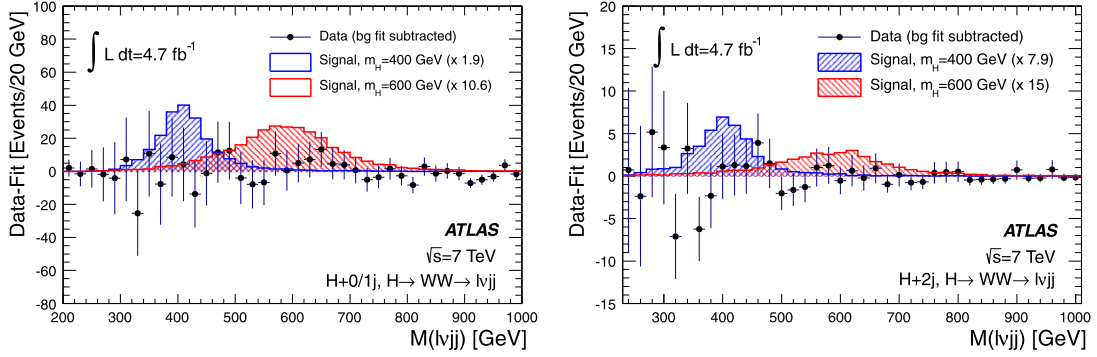


Fig. 6. The difference between data and the fitted background under a no-signal hypothesis, for the (left) $\ell\nu jj + 0/1j$ selection and (right) $\ell\nu jj + 2j$ selection, both summed over lepton flavours. The expected contribution from SM Higgs boson decays is also shown for $m_H = 400$ GeV and $m_H = 600$ GeV, multiplied by a factor equal to the ratio of 95% CL limit on its production to the SM prediction. Uncertainties on the signal normalization and the background shape are not shown in the plots but are taken into account in the limit setting.

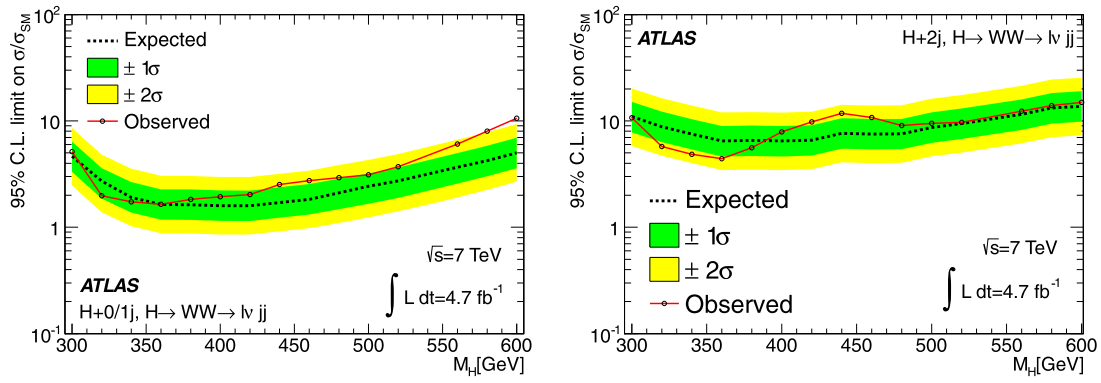


Fig. 7. The expected and observed 95% CL upper limits on the Higgs boson production cross section divided by the SM prediction. The left figure shows the combination of $H + 0j$ with $H + 1j$ and the right figure shows the $H + 2j$ limits. For any hypothesized Higgs boson mass, the background contribution used in the calculation of this limit is obtained from a fit to the $m(\ell\nu jj)$ distribution. The dark (green in the web version) and light (yellow in the web version) bands show the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties on the expected limit.

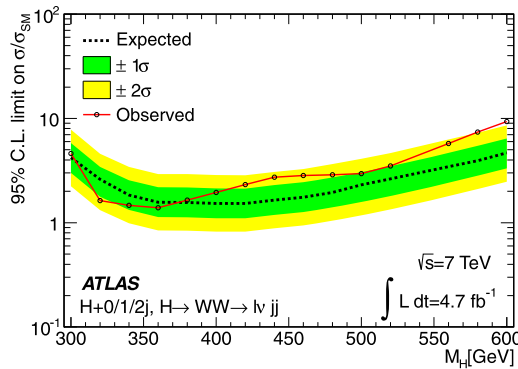


Fig. 8. The expected and observed 95% CL upper limits on the Higgs boson production cross section divided by the SM prediction. This figure shows the combination of the $H + 0j$, $H + 1j$ and $H + 2j$ channels. The background contribution used in the calculation of this limit is obtained from a fit to the $m(\ell\nu jj)$ distribution. The dark (green in the web version) and light (yellow in the web version) bands show the $\pm 1\sigma$ and $\pm 2\sigma$ uncertainties on the expected limit.

large as the one observed in data if there is no signal contribution, where the signal and background are modelled as described in Section 8. The expected p_0 for $H + 0/1j$ if there were a SM Higgs at 400 GeV is 0.091, and the observed value is 0.276. For $H + 2j$, the expected p_0 is 0.369 and the observed is 0.293. The significance is computed as $\sqrt{-2\log\lambda}$ where λ is the likelihood ratio obtained by the fit, and the significance is converted into the probability p_0 using the Gauss error function.

In summary, a search for the SM Higgs boson has been performed in the $H \rightarrow WW \rightarrow \ell\nu jj$ channel using 4.7 fb^{-1} of pp collisions at $\sqrt{s} = 7$ TeV recorded by the ATLAS detector. No

significant excess of events over the expected background has been observed. Exclusion limits on SM Higgs boson production at 95% CL are reported over the Higgs boson mass range of 300–600 GeV.

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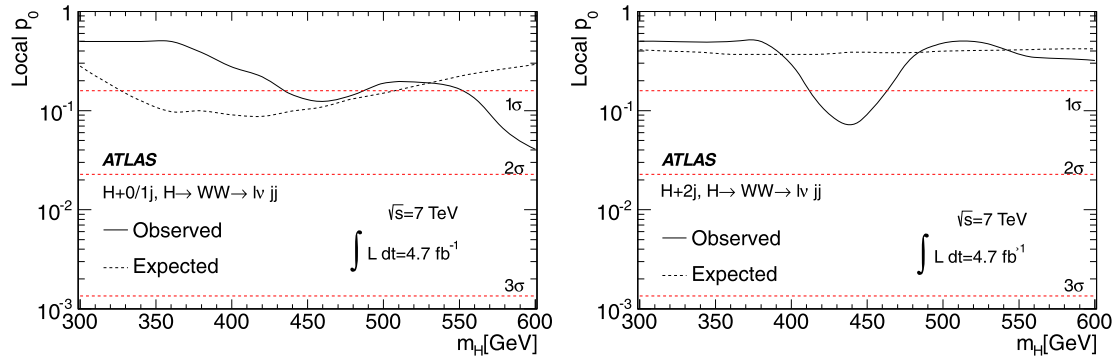


Fig. 9. Local p_0 for the SM Higgs boson search in the $H + 0/1j$ channel (left) and $H + 2j$ channel (right). The dashed line shows the expected p_0 value for a Standard Model Higgs boson as a function of its mass.

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G. Brown⁸¹, H. Brown⁷, P.A. Bruckman de Renstrom³⁸, D. Bruncko^{143b}, R. Bruneliere⁴⁷, S. Brunet⁵⁹, A. Bruni^{19a}, G. Bruni^{19a}, M. Bruschi^{19a}, T. Buanes¹³, Q. Buat⁵⁴, F. Bucci⁴⁸, J. Buchanan¹¹⁷, P. Buchholz¹⁴⁰, R.M. Buckingham¹¹⁷, A.G. Buckley⁴⁵, S.I. Buda^{25a}, I.A. Budagov⁶³, B. Budick¹⁰⁷, V. Büscher⁸⁰, L. Bugge¹¹⁶, O. Bulekov⁹⁵, A.C. Bundock⁷², M. Bunse⁴², T. Buran¹¹⁶, H. Burckhart²⁹, S. Burdin⁷², T. Burgess¹³, S. Burke¹²⁸, E. Busato³³, P. Bussey⁵², C.P. Buszello¹⁶⁵, B. Butler¹⁴², J.M. Butler²¹, C.M. Buttar⁵², J.M. Butterworth⁷⁶, W. Buttinger²⁷, S. Cabrera Urbán¹⁶⁶, D. Caforio^{19a,19b}, O. Cakir^{3a}, P. Calafiura¹⁴, G. Calderini⁷⁷, P. Calfayan⁹⁷, R. Calkins¹⁰⁵, L.P. Caloba^{23a}, R. Caloi^{131a,131b}, D. Calvet³³, S. Calvet³³, R. Camacho Toro³³, P. Camarri^{132a,132b}, D. Cameron¹¹⁶, L.M. Caminada¹⁴, S. Campana²⁹, M. Campanelli⁷⁶, V. Canale^{101a,101b}, F. Canelli^{30,g}, A. Canepa^{158a}, J. Cantero⁷⁹, R. Cantrill⁷⁵, L. Capasso^{101a,101b}, M.D.M. Capeans Garrido²⁹, I. Caprini^{25a}, M. Caprini^{25a}, D. Capriotti⁹⁸, M. Capua^{36a,36b}, R. Caputo⁸⁰, R. Cardarelli^{132a}, T. Carli²⁹, G. Carlino^{101a}, L. Carminati^{88a,88b}, B. Caron⁸⁴, S. Caron¹⁰³, E. Carquin^{31b}, G.D. Carrillo Montoya¹⁷², A.A. Carter⁷⁴, J.R. Carter²⁷, J. Carvalho^{123a,h}, D. Casadei¹⁰⁷, M.P. Casado¹¹, M. Cascella^{121a,121b}, C. Caso^{49a,49b,*}, A.M. Castaneda Hernandez^{172,i}, E. Castaneda-Miranda¹⁷², V. Castillo Gimenez¹⁶⁶, N.F. Castro^{123a}, G. Cataldi^{71a}, P. Catastini⁵⁶, A. Catinaccio²⁹, J.R. Catmore²⁹, A. Cattai²⁹, G. Cattani^{132a,132b}, S. Caughron⁸⁷, P. Cavalleri⁷⁷, D. Cavalli^{88a}, M. Cavalli-Sforza¹¹, V. Cavasinni^{121a,121b}, F. Ceradini^{133a,133b}, A.S. Cerqueira^{23b}, A. Cerri²⁹, L. Cerrito⁷⁴, F. Cerutti⁴⁶, S.A. Cetin^{18b}, A. Chafaq^{134a}, D. Chakraborty¹⁰⁵, I. Chalupkova¹²⁵, K. Chan², B. Chapleau⁸⁴, J.D. Chapman²⁷, J.W. Chapman⁸⁶, E. Chareyre⁷⁷, D.G. Charlton¹⁷, V. Chavda⁸¹, C.A. Chavez Barajas²⁹, S. Cheatham⁸⁴, S. Chekanov⁵, S.V. Chekulaev^{158a}, G.A. Chelkov⁶³, M.A. Chelstowska¹⁰³, C. Chen⁶², H. Chen²⁴, S. Chen^{32c}, X. Chen¹⁷², Y. Chen³⁴, A. Cheplakov⁶³, R. Cherkaoui El Moursli^{134e}, V. Chernyatin²⁴, E. Cheu⁶, S.L. Cheung¹⁵⁷, L. Chevalier¹³⁵, G. Chiefari^{101a,101b}, L. Chikovani^{50a,*}, J.T. Childers²⁹, A. Chilingarov⁷⁰, G. Chiodini^{71a}, A.S. Chisholm¹⁷, R.T. Chislett⁷⁶, A. Chitan^{25a}, M.V. Chizhov⁶³, G. Choudalakis³⁰, S. Chouridou¹³⁶, I.A. Christidi⁷⁶, A. Christov⁴⁷, D. Chromek-Burckhart²⁹, M.L. Chu¹⁵⁰, J. Chudoba¹²⁴, G. Ciapetti^{131a,131b}, A.K. Ciftci^{3a}, R. Ciftci^{3a}, D. Cinca³³, V. Cindro⁷³, C. Ciocca^{19a,19b}, A. Ciochio¹⁴, M. Cirilli⁸⁶, P. Cirkovic^{12b}, M. Citterio^{88a}, M. Ciubancan^{25a}, A. Clark⁴⁸, P.J. Clark⁴⁵, R.N. Clarke¹⁴, W. Cleland¹²², J.C. Clemens⁸², B. Clement⁵⁴, C. Clement^{145a,145b}, Y. Coadou⁸², M. Cobl^{163a,163c}, A. Coccaro¹³⁷, J. Cochran⁶², J.G. Cogan¹⁴², J. Coggeshall¹⁶⁴, E. Cogneras¹⁷⁷, J. Colas⁴, S. Cole¹⁰⁵, A.P. Colijn¹⁰⁴, N.J. Collins¹⁷, C. Collins-Tooth⁵², J. Collot⁵⁴, T. Colombo^{118a,118b}, G. Colon⁸³, P. Conde Muiño^{123a}, E. Coniavitis¹¹⁷, M.C. Conidi¹¹, S.M. Consonni^{88a,88b}, V. Consorti⁴⁷, S. Constantinescu^{25a}, C. Conta^{118a,118b}, G. Conti⁵⁶, F. Conventi^{101a,j}, M. Cooke¹⁴, B.D. Cooper⁷⁶, A.M. Cooper-Sarkar¹¹⁷, K. Copic¹⁴, T. Cornelissen¹⁷⁴, M. Corradi^{19a}, F. Corriveau^{84,k}, A. Cortes-Gonzalez¹⁶⁴, G. Cortiana⁹⁸, G. Costa^{88a}, M.J. Costa¹⁶⁶, D. Costanzo¹³⁸, T. Costin³⁰, D. Côté²⁹, L. Courneyea¹⁶⁸, G. Cowan⁷⁵, C. Cowden²⁷, B.E. Cox⁸¹, K. Cranmer¹⁰⁷, F. Crescioli^{121a,121b}, M. Cristinziani²⁰, G. Crosetti^{36a,36b}, S. Crépe-Renaudin⁵⁴, C.-M. Cuciuc^{25a}, C. Cuenca Almenar¹⁷⁵, T. Cuhadar Donszelmann¹³⁸, M. Curatolo⁴⁶, C.J. Curtis¹⁷, C. Cuthbert¹⁴⁹, P. Cwetanski⁵⁹, H. Czirr¹⁴⁰, P. Czodrowski⁴³, Z. Czyzula¹⁷⁵, S. D'Auria⁵², M. D'Onofrio⁷², A. D'Orazio^{131a,131b}, M.J. Da Cunha Sargedas De Sousa^{123a}, C. Da Via⁸¹, W. Dabrowski³⁷, A. Dafinca¹¹⁷, T. Dai⁸⁶, C. Dallapiccola⁸³, M. Dam³⁵, M. Dameri^{49a,49b}, D.S. Damiani¹³⁶, H.O. Danielsson²⁹, V. Dao⁴⁸, G. Darbo^{49a}, G.L. Darlea^{25b}, J.A. Dassoulas⁴¹, W. Davey²⁰, T. Davidek¹²⁵, N. Davidson⁸⁵, R. Davidson⁷⁰, E. Davies^{117,c}, M. Davies⁹², O. Davignon⁷⁷, A.R. Davison⁷⁶, Y. Davygora^{57a}, E. Dawe¹⁴¹, I. Dawson¹³⁸, R.K. Daya-Ishmukhametova²², K. De⁷, R. de Asmundis^{101a}, S. De Castro^{19a,19b}, S. De Cecco⁷⁷, J. de Graat⁹⁷, N. De Groot¹⁰³, P. de Jong¹⁰⁴, C. De La Taille¹¹⁴, H. De la Torre⁷⁹, F. De Lorenzi⁶², L. de Mora⁷⁰, L. De Nooij¹⁰⁴, D. De Pedis^{131a}, A. De Salvo^{131a}, U. De Sanctis^{163a,163c}, A. De Santo¹⁴⁸, J.B. De Vivie De Regie¹¹⁴, G. De Zorzi^{131a,131b}, W.J. Dearnaley⁷⁰, R. Debbe²⁴, C. Debenedetti⁴⁵, B. Dechenaux⁵⁴, D.V. Dedovich⁶³, J. Degenhardt¹¹⁹, C. Del Papa^{163a,163c}, J. Del Peso⁷⁹, T. Del Prete^{121a,121b}, T. Delemontex⁵⁴, M. Deliyergiyev⁷³, A. Dell'Acqua²⁹, L. Dell'Asta²¹, M. Della Pietra^{101a,j}, D. della Volpe^{101a,101b}, M. Delmastro⁴, P.A. Delsart⁵⁴, C. Deluca¹⁰⁴, S. Demers¹⁷⁵, M. Demichev⁶³, B. Demirkoz^{11,l}, J. Deng¹⁶², S.P. Denisov¹²⁷, D. Derendarz³⁸, J.E. Derkaoui^{134d}, F. Derue⁷⁷, P. Dervan⁷², K. Desch²⁰, E. Devetak¹⁴⁷, P.O. Deviveiros¹⁰⁴, A. Dewhurst¹²⁸, B. DeWilde¹⁴⁷, S. Dhaliwal¹⁵⁷, R. Dhullipudi^{24,m}, A. Di Ciaccio^{132a,132b}, L. Di Ciaccio⁴, A. Di Girolamo²⁹, B. Di Girolamo²⁹, S. Di Luise^{133a,133b}, A. Di Mattia¹⁷², B. Di Micco²⁹, R. Di Nardo⁴⁶, A. Di Simone^{132a,132b}, R. Di Sipio^{19a,19b}, M.A. Diaz^{31a}, E.B. Diehl⁸⁶, J. Dietrich⁴¹, T.A. Dietzsch^{57a},

S. Diglio⁸⁵, K. Dindar Yagci³⁹, J. Dingfelder²⁰, F. Dinut^{25a}, C. Dionisi^{131a,131b}, P. Dita^{25a}, S. Dita^{25a}, F. Dittus²⁹, F. Djama⁸², T. Djobava^{50b}, M.A.B. do Vale^{23c}, A. Do Valle Wemans^{123a,n}, T.K.O. Doan⁴, M. Dobbs⁸⁴, R. Dobinson^{29,*}, D. Dobos²⁹, E. Dobson^{29,o}, J. Dodd³⁴, C. Doglioni⁴⁸, T. Doherty⁵², Y. Doi^{64,*}, J. Dolejsi¹²⁵, I. Dolenc⁷³, Z. Dolezal¹²⁵, B.A. Dolgoshein^{95,*}, T. Dohmae¹⁵⁴, M. Donadelli^{23d}, J. Donini³³, J. Dopke²⁹, A. Doria^{101a}, A. Dos Anjos¹⁷², A. Dotti^{121a,121b}, M.T. Dova⁶⁹, A.D. Doxiadis¹⁰⁴, A.T. Doyle⁵², M. Dris⁹, J. Dubbert⁹⁸, S. Dube¹⁴, E. Duchovni¹⁷¹, G. Duckeck⁹⁷, A. Dudarev²⁹, F. Dudziak⁶², M. Dührssen²⁹, I.P. Duerdoth⁸¹, L. Duflot¹¹⁴, M.-A. Dufour⁸⁴, L. Duguid⁷⁵, M. Dunford²⁹, H. Duran Yildiz^{3a}, R. Duxfield¹³⁸, M. Dwuznik³⁷, F. Dydak²⁹, M. Düren⁵¹, J. Ebke⁹⁷, S. Eckweiler⁸⁰, K. Edmonds⁸⁰, W. Edson¹, C.A. Edwards⁷⁵, N.C. Edwards⁵², W. Ehrenfeld⁴¹, T. Eifert¹⁴², G. Eigen¹³, K. Einsweiler¹⁴, E. Eisenhandler⁷⁴, T. Ekelof¹⁶⁵, M. El Kacimi^{134c}, M. Ellert¹⁶⁵, S. Elles⁴, F. Ellinghaus⁸⁰, K. Ellis⁷⁴, N. Ellis²⁹, J. Elmsheuser⁹⁷, M. Elsing²⁹, D. Emeliyanov¹²⁸, R. Engelmann¹⁴⁷, A. Engl⁹⁷, B. Epp⁶⁰, J. Erdmann⁵³, A. Ereditato¹⁶, D. Eriksson^{145a}, J. Ernst¹, M. Ernst²⁴, J. Ernwein¹³⁵, D. Errede¹⁶⁴, S. Errede¹⁶⁴, E. Ertel⁸⁰, M. Escalier¹¹⁴, H. Esch⁴², C. Escobar¹²², X. Espinal Curull¹¹, B. Esposito⁴⁶, F. Etienne⁸², A.I. Etienvre¹³⁵, E. Etzion¹⁵², D. Evangelakou⁵³, H. Evans⁵⁹, L. Fabbri^{19a,19b}, C. Fabre²⁹, R.M. Fakhruddinov¹²⁷, S. Falciano^{131a}, Y. Fang¹⁷², M. Fanti^{88a,88b}, A. Farbin⁷, A. Farilla^{133a}, J. Farley¹⁴⁷, T. Farooque¹⁵⁷, S. Farrell¹⁶², S.M. Farrington¹⁶⁹, P. Farthouat²⁹, P. Fassnacht²⁹, D. Fassouliotis⁸, B. Fatholahzadeh¹⁵⁷, A. Favareto^{88a,88b}, L. Fayard¹¹⁴, S. Fazio^{36a,36b}, R. Febbraro³³, P. Federic^{143a}, O.L. Fedin¹²⁰, W. Fedorko⁸⁷, M. Fehling-Kaschek⁴⁷, L. Feligioni⁸², D. Fellmann⁵, C. Feng^{32d}, E.J. Feng⁵, A.B. Fenyuk¹²⁷, J. Ferencei^{143b}, W. Fernando⁵, S. Ferrag⁵², J. Ferrando⁵², V. Ferrara⁴¹, A. Ferrari¹⁶⁵, P. Ferrari¹⁰⁴, R. Ferrari^{118a}, D.E. Ferreira de Lima⁵², A. Ferrer¹⁶⁶, D. Ferrere⁴⁸, C. Ferretti⁸⁶, A. Ferretto Parodi^{49a,49b}, M. Fiascaris³⁰, F. Fiedler⁸⁰, A. Filipčič⁷³, F. Filthaut¹⁰³, M. Fincke-Keeler¹⁶⁸, M.C.N. Fiolhais^{123a,h}, L. Fiorini¹⁶⁶, A. Firan³⁹, G. Fischer⁴¹, M.J. Fisher¹⁰⁸, M. Flechl⁴⁷, I. Fleck¹⁴⁰, J. Fleckner⁸⁰, P. Fleischmann¹⁷³, S. Fleischmann¹⁷⁴, T. Flick¹⁷⁴, A. Floderus⁷⁸, L.R. Flores Castillo¹⁷², M.J. Flowerdew⁹⁸, T. Fonseca Martin¹⁶, A. Formica¹³⁵, A. Forti⁸¹, D. Fortin^{158a}, D. Fournier¹¹⁴, H. Fox⁷⁰, P. Francavilla¹¹, M. Franchini^{19a,19b}, S. Franchino^{118a,118b}, D. Francis²⁹, T. Frank¹⁷¹, S. Franz²⁹, M. Fraternali^{118a,118b}, S. Fratina¹¹⁹, S.T. French²⁷, C. Friedrich⁴¹, F. Friedrich⁴³, R. Froeschl²⁹, D. Froidevaux²⁹, J.A. Frost²⁷, C. Fukunaga¹⁵⁵, E. Fullana Torregrosa²⁹, B.G. Fulsom¹⁴², J. Fuster¹⁶⁶, C. Gabaldon²⁹, O. Gabizon¹⁷¹, T. Gadfort²⁴, S. Gadomski⁴⁸, G. Gagliardi^{49a,49b}, P. Gagnon⁵⁹, C. Galea⁹⁷, E.J. Gallas¹¹⁷, V. Gallo¹⁶, B.J. Gallop¹²⁸, P. Gallus¹²⁴, K.K. Gan¹⁰⁸, Y.S. Gao^{142,e}, A. Gaponenko¹⁴, F. Garbersson¹⁷⁵, M. Garcia-Sciveres¹⁴, C. García¹⁶⁶, J.E. García Navarro¹⁶⁶, R.W. Gardner³⁰, N. Garelli²⁹, H. Garitaonandia¹⁰⁴, V. Garonne²⁹, J. Garvey¹⁷, C. Gatti⁴⁶, G. Gaudio^{118a}, B. Gaur¹⁴⁰, L. Gauthier¹³⁵, P. Gauzzi^{131a,131b}, I.L. Gavrilenko⁹³, C. Gay¹⁶⁷, G. Gaycken²⁰, E.N. Gazis⁹, P. Ge^{32d}, Z. Gecse¹⁶⁷, C.N.P. Gee¹²⁸, D.A.A. Geerts¹⁰⁴, Ch. Geich-Gimbel²⁰, K. Gellerstedt^{145a,145b}, C. Gemme^{49a}, A. Gemmell⁵², M.H. Genest⁵⁴, S. Gentile^{131a,131b}, M. George⁵³, S. George⁷⁵, P. Gerlach¹⁷⁴, A. Gershon¹⁵², C. Geweniger^{57a}, H. Ghazlane^{134b}, N. Ghodbane³³, B. Giacobbe^{19a}, S. Giagu^{131a,131b}, V. Giakoumopoulou⁸, V. Giangiobbe¹¹, F. Gianotti²⁹, B. Gibbard²⁴, A. Gibson¹⁵⁷, S.M. Gibson²⁹, D. Gillberg²⁸, A.R. Gillman¹²⁸, D.M. Gingrich^{2,d}, J. Ginzburg¹⁵², N. Giokaris⁸, M.P. Giordani^{163c}, R. Giordano^{101a,101b}, F.M. Giorgi¹⁵, P. Giovannini⁹⁸, P.F. Giraud¹³⁵, D. Giugni^{88a}, M. Giunta⁹², P. Giusti^{19a}, B.K. Gjelsten¹¹⁶, L.K. Gladilin⁹⁶, C. Glasman⁷⁹, J. Glatzer⁴⁷, A. Glazov⁴¹, K.W. Glitza¹⁷⁴, G.L. Glonti⁶³, J.R. Goddard⁷⁴, J. Godfrey¹⁴¹, J. Godlewski²⁹, M. Goebel⁴¹, T. Göpfert⁴³, C. Goeringer⁸⁰, C. Gössling⁴², S. Goldfarb⁸⁶, T. Golling¹⁷⁵, A. Gomes^{123a,b}, L.S. Gomez Fajardo⁴¹, R. Gonçalves⁷⁵, J. Goncalves Pinto Firmino Da Costa⁴¹, L. Gonella²⁰, S. Gonzalez¹⁷², S. González de la Hoz¹⁶⁶, G. Gonzalez Parra¹¹, M.L. Gonzalez Silva²⁶, S. Gonzalez-Sevilla⁴⁸, J.J. Goodson¹⁴⁷, L. Goossens²⁹, P.A. Gorbounov⁹⁴, H.A. Gordon²⁴, I. Gorelov¹⁰², G. Gorfine¹⁷⁴, B. Gorini²⁹, E. Gorini^{71a,71b}, A. Gorišek⁷³, E. Gornicki³⁸, B. Gosdzik⁴¹, A.T. Goshaw⁵, M. Gosselink¹⁰⁴, M.I. Gostkin⁶³, I. Gough Eschrich¹⁶², M. Gouighri^{134a}, D. Goujdami^{134c}, M.P. Goulette⁴⁸, A.G. Goussiou¹³⁷, C. Goy⁴, S. Gozpinar²², I. Grabowska-Bold³⁷, P. Grafström^{19a,19b}, K.-J. Grahn⁴¹, F. Grancagnolo^{71a}, S. Grancagnolo¹⁵, V. Grassi¹⁴⁷, V. Gratchev¹²⁰, N. Grau³⁴, H.M. Gray²⁹, J.A. Gray¹⁴⁷, E. Graziani^{133a}, O.G. Grebenyuk¹²⁰, T. Greenshaw⁷², Z.D. Greenwood^{24,m}, K. Gregersen³⁵, I.M. Gregor⁴¹, P. Grenier¹⁴², J. Griffiths¹³⁷, N. Grigalashvili⁶³, A.A. Grillo¹³⁶, S. Grinstein¹¹, Y.V. Grishkevich⁹⁶, J.-F. Grivaz¹¹⁴, E. Gross¹⁷¹, J. Grosse-Knetter⁵³, J. Groth-Jensen¹⁷¹, K. Grybel¹⁴⁰, D. Guest¹⁷⁵, C. Guicheney³³, S. Guindon⁵³, U. Gul⁵², H. Guler^{84,p}, J. Gunther¹²⁴, B. Guo¹⁵⁷, J. Guo³⁴, P. Gutierrez¹¹⁰,

N. Guttman¹⁵², O. Gutzwiller¹⁷², C. Guyot¹³⁵, C. Gwenlan¹¹⁷, C.B. Gwilliam⁷², A. Haas¹⁴², S. Haas²⁹,
 C. Haber¹⁴, H.K. Hadavand³⁹, D.R. Hadley¹⁷, P. Haefner²⁰, F. Hahn²⁹, S. Haider²⁹, Z. Hajduk³⁸,
 H. Hakobyan¹⁷⁶, D. Hall¹¹⁷, J. Haller⁵³, K. Hamacher¹⁷⁴, P. Hamal¹¹², M. Hamer⁵³, A. Hamilton^{144b,q},
 S. Hamilton¹⁶⁰, L. Han^{32b}, K. Hanagaki¹¹⁵, K. Hanawa¹⁵⁹, M. Hance¹⁴, C. Handel⁸⁰, P. Hanke^{57a},
 J.R. Hansen³⁵, J.B. Hansen³⁵, J.D. Hansen³⁵, P.H. Hansen³⁵, P. Hansson¹⁴², K. Hara¹⁵⁹, G.A. Hare¹³⁶,
 T. Harenberg¹⁷⁴, S. Harkusha⁸⁹, D. Harper⁸⁶, R.D. Harrington⁴⁵, O.M. Harris¹³⁷, J. Hartert⁴⁷,
 F. Hartjes¹⁰⁴, T. Haruyama⁶⁴, A. Harvey⁵⁵, S. Hasegawa¹⁰⁰, Y. Hasegawa¹³⁹, S. Hassani¹³⁵, S. Haug¹⁶,
 M. Hauschild²⁹, R. Hauser⁸⁷, M. Havranek²⁰, C.M. Hawkes¹⁷, R.J. Hawkins²⁹, A.D. Hawkins⁷⁸,
 D. Hawkins¹⁶², T. Hayakawa⁶⁵, T. Hayashi¹⁵⁹, D. Hayden⁷⁵, C.P. Hays¹¹⁷, H.S. Hayward⁷²,
 S.J. Haywood¹²⁸, M. He^{32d}, S.J. Head¹⁷, V. Hedberg⁷⁸, L. Heelan⁷, S. Heim⁸⁷, B. Heinemann¹⁴,
 S. Heisterkamp³⁵, L. Helary²¹, C. Heller⁹⁷, M. Heller²⁹, S. Hellman^{145a,145b}, D. Hellmich²⁰, C. Helsen¹¹,
 R.C.W. Henderson⁷⁰, M. Henke^{57a}, A. Henrichs⁵³, A.M. Henriques Correia²⁹, S. Henrot-Versille¹¹⁴,
 C. Hensel⁵³, T. Henß¹⁷⁴, C.M. Hernandez⁷, Y. Hernández Jiménez¹⁶⁶, R. Herrberg¹⁵, G. Herten⁴⁷,
 R. Hertenberger⁹⁷, L. Hervas²⁹, G.G. Hesketh⁷⁶, N.P. Hessey¹⁰⁴, E. Higón-Rodríguez¹⁶⁶, J.C. Hill²⁷,
 K.H. Hiller⁴¹, S. Hillert²⁰, S.J. Hillier¹⁷, I. Hinchliffe¹⁴, E. Hines¹¹⁹, M. Hirose¹¹⁵, F. Hirsch⁴²,
 D. Hirschbuehl¹⁷⁴, J. Hobbs¹⁴⁷, N. Hod¹⁵², M.C. Hodgkinson¹³⁸, P. Hodgson¹³⁸, A. Hoecker²⁹,
 M.R. Hoferkamp¹⁰², J. Hoffman³⁹, D. Hoffmann⁸², M. Hohlfeld⁸⁰, M. Holder¹⁴⁰, S.O. Holmgren^{145a},
 T. Holy¹²⁶, J.L. Holzbauer⁸⁷, T.M. Hong¹¹⁹, L. Hooft van Huysduynen¹⁰⁷, C. Horn¹⁴², S. Horner⁴⁷,
 J.-Y. Hostachy⁵⁴, S. Hou¹⁵⁰, A. Hoummada^{134a}, J. Howard¹¹⁷, J. Howarth⁸¹, I. Hristova¹⁵, J. Hrivnac¹¹⁴,
 T. Hryn'ova⁴, P.J. Hsu⁸⁰, S.-C. Hsu¹⁴, Z. Hubacek¹²⁶, F. Hubaut⁸², F. Huegging²⁰, A. Huettmann⁴¹,
 T.B. Huffman¹¹⁷, E.W. Hughes³⁴, G. Hughes⁷⁰, M. Huhtinen²⁹, M. Hurwitz¹⁴, U. Husemann⁴¹,
 N. Huseynov^{63,r}, J. Huston⁸⁷, J. Huth⁵⁶, G. Iacobucci⁴⁸, G. Iakovidis⁹, M. Ibbotson⁸¹, I. Ibragimov¹⁴⁰,
 L. Iconomidou-Fayard¹¹⁴, J. Idarraga¹¹⁴, P. Iengo^{101a}, O. Igonkina¹⁰⁴, Y. Ikegami⁶⁴, M. Ikeno⁶⁴,
 D. Iliadis¹⁵³, N. Ilic¹⁵⁷, T. Ince²⁰, J. Inigo-Golfín²⁹, P. Ioannou⁸, M. Iodice^{133a}, K. Iordanidou⁸,
 V. Ippolito^{131a,131b}, A. Irlles Quiles¹⁶⁶, C. Isaksson¹⁶⁵, M. Ishino⁶⁶, M. Ishitsuka¹⁵⁶, R. Ishmukhametov³⁹,
 C. Issever¹¹⁷, S. Istin^{18a}, A.V. Ivashin¹²⁷, W. Iwanski³⁸, H. Iwasaki⁶⁴, J.M. Izen⁴⁰, V. Izzo^{101a},
 B. Jackson¹¹⁹, J.N. Jackson⁷², P. Jackson¹⁴², M.R. Jaekel²⁹, V. Jain⁵⁹, K. Jakobs⁴⁷, S. Jakobsen³⁵,
 T. Jakoubek¹²⁴, J. Jakubek¹²⁶, D.K. Jana¹¹⁰, E. Jansen⁷⁶, H. Jansen²⁹, A. Jantsch⁹⁸, M. Janus⁴⁷,
 G. Jarlskog⁷⁸, L. Jeanty⁵⁶, I. Jen-La Plante³⁰, D. Jennens⁸⁵, P. Jenni²⁹, P. Jež³⁵, S. Jézéquel⁴, M.K. Jha^{19a},
 H. Ji¹⁷², W. Ji⁸⁰, J. Jia¹⁴⁷, Y. Jiang^{32b}, M. Jimenez Belenguer⁴¹, S. Jin^{32a}, O. Jinnouchi¹⁵⁶,
 M.D. Joergensen³⁵, D. Joffe³⁹, M. Johansen^{145a,145b}, K.E. Johansson^{145a}, P. Johansson¹³⁸, S. Johnert⁴¹,
 K.A. Johns⁶, K. Jon-And^{145a,145b}, G. Jones¹⁶⁹, R.W.L. Jones⁷⁰, T.J. Jones⁷², C. Joram²⁹, P.M. Jorge^{123a},
 K.D. Joshi⁸¹, J. Jovicevic¹⁴⁶, T. Jovin^{12b}, X. Ju¹⁷², C.A. Jung⁴², R.M. Jungst²⁹, V. Juranek¹²⁴, P. Jussel⁶⁰,
 A. Juste Rozas¹¹, S. Kabana¹⁶, M. Kaci¹⁶⁶, A. Kaczmarska³⁸, P. Kadlecik³⁵, M. Kado¹¹⁴, H. Kagan¹⁰⁸,
 M. Kagan⁵⁶, E. Kajomovitz¹⁵¹, S. Kalinin¹⁷⁴, L.V. Kalinovskaya⁶³, S. Kama³⁹, N. Kanaya¹⁵⁴, M. Kaneda²⁹,
 S. Kaneti²⁷, T. Kanno¹⁵⁶, V.A. Kantserov⁹⁵, J. Kanzaki⁶⁴, B. Kaplan¹⁷⁵, A. Kapliy³⁰, J. Kaplon²⁹, D. Kar⁵²,
 M. Karagounis²⁰, K. Karakostas⁹, M. Karnevskiy⁴¹, V. Kartvelishvili⁷⁰, A.N. Karyukhin¹²⁷, L. Kashif¹⁷²,
 G. Kasieczka^{57b}, R.D. Kass¹⁰⁸, A. Kastanas¹³, M. Kataoka⁴, Y. Kataoka¹⁵⁴, E. Katsoufis⁹, J. Katzy⁴¹,
 V. Kaushik⁶, K. Kawagoe⁶⁸, T. Kawamoto¹⁵⁴, G. Kawamura⁸⁰, M.S. Kayl¹⁰⁴, V.A. Kazanin¹⁰⁶,
 M.Y. Kazarinov⁶³, R. Keeler¹⁶⁸, R. Kehoe³⁹, M. Keil⁵³, G.D. Kekelidze⁶³, J.S. Keller¹³⁷, M. Kenyon⁵²,
 O. Kepka¹²⁴, N. Kerschen²⁹, B.P. Kerševan⁷³, S. Kersten¹⁷⁴, K. Kessoku¹⁵⁴, J. Keung¹⁵⁷, F. Khalil-zada¹⁰,
 H. Khandanyan¹⁶⁴, A. Khanov¹¹¹, D. Kharchenko⁶³, A. Khodinov⁹⁵, A. Khomich^{57a}, T.J. Khoo²⁷,
 G. Khoraiuli²⁰, A. Khoroshilov¹⁷⁴, V. Khovanskiy⁹⁴, E. Khramov⁶³, J. Khubua^{50b}, H. Kim^{145a,145b},
 S.H. Kim¹⁵⁹, N. Kimura¹⁷⁰, O. Kind¹⁵, B.T. King⁷², M. King⁶⁵, R.S.B. King¹¹⁷, J. Kirk¹²⁸, A.E. Kiryunin⁹⁸,
 T. Kishimoto⁶⁵, D. Kisielewska³⁷, T. Kitamura⁶⁵, T. Kittelmann¹²², E. Kladiva^{143b}, M. Klein⁷², U. Klein⁷²,
 K. Kleinknecht⁸⁰, M. Klemetti⁸⁴, A. Klier¹⁷¹, P. Klimek^{145a,145b}, A. Klimentov²⁴, R. Klingenberg⁴²,
 J.A. Klinger⁸¹, E.B. Klinkby³⁵, T. Klioutchnikova²⁹, P.F. Klok¹⁰³, S. Klous¹⁰⁴, E.-E. Kluge^{57a}, T. Kluge⁷²,
 P. Kluit¹⁰⁴, S. Kluth⁹⁸, N.S. Knecht¹⁵⁷, E. Kneringer⁶⁰, E.B.F.G. Knoop⁸², A. Knue⁵³, B.R. Ko⁴⁴,
 T. Kobayashi¹⁵⁴, M. Kobel⁴³, M. Kocian¹⁴², P. Kodys¹²⁵, K. Köneke²⁹, A.C. König¹⁰³, S. Koenig⁸⁰,
 L. Köpke⁸⁰, F. Koetsveld¹⁰³, P. Koevesarki²⁰, T. Koffas²⁸, E. Koffeman¹⁰⁴, L.A. Kogan¹¹⁷, S. Kohlmann¹⁷⁴,
 F. Kohn⁵³, Z. Kohout¹²⁶, T. Kohriki⁶⁴, T. Koi¹⁴², G.M. Kolachev^{106,*}, H. Kolanoski¹⁵, V. Kolesnikov⁶³,
 I. Koletsou^{88a}, J. Koll⁸⁷, M. Kollefrath⁴⁷, A.A. Komar⁹³, Y. Komori¹⁵⁴, T. Kondo⁶⁴, T. Kono^{41,s},

A.I. Kononov⁴⁷, R. Konoplich^{107,t}, N. Konstantinidis⁷⁶, S. Koperny³⁷, K. Korcyl³⁸, K. Kordas¹⁵³,
 A. Korn¹¹⁷, A. Korol¹⁰⁶, I. Korolkov¹¹, E.V. Korolkova¹³⁸, V.A. Korotkov¹²⁷, O. Kortner⁹⁸, S. Kortner⁹⁸,
 V.V. Kostyukhin²⁰, S. Kotov⁹⁸, V.M. Kotov⁶³, A. Kotwal⁴⁴, C. Kourkouvelis⁸, V. Kouskoura¹⁵³,
 A. Koutsman^{158a}, R. Kowalewski¹⁶⁸, T.Z. Kowalski³⁷, W. Kozanecki¹³⁵, A.S. Kozhin¹²⁷, V. Kral¹²⁶,
 V.A. Kramarenko⁹⁶, G. Kramberger⁷³, M.W. Krasny⁷⁷, A. Krasznahorkay¹⁰⁷, J.K. Kraus²⁰, S. Kreiss¹⁰⁷,
 F. Krejci¹²⁶, J. Kretschmar⁷², N. Krieger⁵³, P. Krieger¹⁵⁷, K. Kroeninger⁵³, H. Kroha⁹⁸, J. Kroll¹¹⁹,
 J. Kroseberg²⁰, J. Krstic^{12a}, U. Kruchonak⁶³, H. Krüger²⁰, T. Kruker¹⁶, N. Krumnack⁶²,
 Z.V. Krumshteyn⁶³, T. Kubota⁸⁵, S. Kuday^{3a}, S. Kuehn⁴⁷, A. Kugel^{57c}, T. Kuhl⁴¹, D. Kuhn⁶⁰,
 V. Kukhtin⁶³, Y. Kulchitsky⁸⁹, S. Kuleshov^{31b}, C. Kummer⁹⁷, M. Kuna⁷⁷, J. Kunkle¹¹⁹, A. Kupco¹²⁴,
 H. Kurashige⁶⁵, M. Kurata¹⁵⁹, Y.A. Kurochkin⁸⁹, V. Kus¹²⁴, E.S. Kuwertz¹⁴⁶, M. Kuze¹⁵⁶, J. Kvita¹⁴¹,
 R. Kwee¹⁵, A. La Rosa⁴⁸, L. La Rotonda^{36a,36b}, L. Labarga⁷⁹, J. Labbe⁴, S. Lablak^{134a}, C. Lacasta¹⁶⁶,
 F. Lacava^{131a,131b}, H. Lacker¹⁵, D. Lacour⁷⁷, V.R. Lacuesta¹⁶⁶, E. Ladygin⁶³, R. Lafaye⁴, B. Laforge⁷⁷,
 T. Lagouri⁷⁹, S. Lai⁴⁷, E. Laisne⁵⁴, M. Lamanna²⁹, L. Lambourne⁷⁶, C.L. Lampen⁶, W. Lampl⁶,
 E. Lancon¹³⁵, U. Landgraf⁴⁷, M.P.J. Landon⁷⁴, J.L. Lane⁸¹, V.S. Lang^{57a}, C. Lange⁴¹, A.J. Lankford¹⁶²,
 F. Lanni²⁴, K. Lantzsch¹⁷⁴, S. Laplace⁷⁷, C. Lapoire²⁰, J.F. Laporte¹³⁵, T. Lari^{88a}, A. Larner¹¹⁷,
 M. Lassnig²⁹, P. Laurelli⁴⁶, V. Lavorini^{36a,36b}, W. Lavrijsen¹⁴, P. Laycock⁷², O. Le Dortz⁷⁷,
 E. Le Guirriec⁸², C. Le Maner¹⁵⁷, E. Le Menedeu¹¹, T. LeCompte⁵, F. Ledroit-Guillon⁵⁴, H. Lee¹⁰⁴,
 J.S.H. Lee¹¹⁵, S.C. Lee¹⁵⁰, L. Lee¹⁷⁵, M. Lefebvre¹⁶⁸, M. Legendre¹³⁵, F. Legger⁹⁷, C. Leggett¹⁴,
 M. Lehmacher²⁰, G. Lehmann Miotto²⁹, X. Lei⁶, M.A.L. Leite^{23d}, R. Leitner¹²⁵, D. Lellouch¹⁷¹,
 B. Lemmer⁵³, V. Lendermann^{57a}, K.J.C. Leney^{144b}, T. Lenz¹⁰⁴, G. Lenzen¹⁷⁴, B. Lenzi²⁹, K. Leonhardt⁴³,
 S. Leontsinis⁹, F. Lepold^{57a}, C. Leroy⁹², J.-R. Lessard¹⁶⁸, C.G. Lester²⁷, C.M. Lester¹¹⁹, J. Levêque⁴,
 D. Levin⁸⁶, L.J. Levinson¹⁷¹, A. Lewis¹¹⁷, G.H. Lewis¹⁰⁷, A.M. Leyko²⁰, M. Leyton¹⁵, B. Li⁸², H. Li^{172,u},
 S. Li^{32b,v}, X. Li⁸⁶, Z. Liang^{117,w}, H. Liao³³, B. Liberti^{132a}, P. Lichard²⁹, M. Lichtnecker⁹⁷, K. Lie¹⁶⁴,
 W. Liebig¹³, C. Limbach²⁰, A. Limosani⁸⁵, M. Limper⁶¹, S.C. Lin^{150,x}, F. Linde¹⁰⁴, J.T. Linnemann⁸⁷,
 E. Lipeles¹¹⁹, A. Lipniacka¹³, T.M. Liss¹⁶⁴, D. Lissauer²⁴, A. Lister⁴⁸, A.M. Litke¹³⁶, C. Liu²⁸, D. Liu¹⁵⁰,
 H. Liu⁸⁶, J.B. Liu⁸⁶, L. Liu⁸⁶, M. Liu^{32b}, Y. Liu^{32b}, M. Livan^{118a,118b}, S.S.A. Livermore¹¹⁷, A. Lleres⁵⁴,
 J. Llorente Merino⁷⁹, S.L. Lloyd⁷⁴, E. Lobodzinska⁴¹, P. Loch⁶, W.S. Lockman¹³⁶, T. Loddenkoetter²⁰,
 F.K. Loebinger⁸¹, A. Loginov¹⁷⁵, C.W. Loh¹⁶⁷, T. Lohse¹⁵, K. Lohwasser⁴⁷, M. Lokajicek¹²⁴,
 V.P. Lombardo⁴, R.E. Long⁷⁰, L. Lopes^{123a}, D. Lopez Mateos⁵⁶, J. Lorenz⁹⁷, N. Lorenzo Martinez¹¹⁴,
 M. Losada¹⁶¹, P. Loscutoff¹⁴, F. Lo Sterzo^{131a,131b}, M.J. Losty^{158a}, X. Lou⁴⁰, A. Lounis¹¹⁴,
 K.F. Loureiro¹⁶¹, J. Love²¹, P.A. Love⁷⁰, A.J. Lowe^{142,e}, F. Lu^{32a}, H.J. Lubatti¹³⁷, C. Luci^{131a,131b},
 A. Lucotte⁵⁴, A. Ludwig⁴³, D. Ludwig⁴¹, I. Ludwig⁴⁷, J. Ludwig⁴⁷, F. Luehring⁵⁹, G. Luijckx¹⁰⁴,
 W. Lukas⁶⁰, D. Lumb⁴⁷, L. Luminari^{131a}, E. Lund¹¹⁶, B. Lund-Jensen¹⁴⁶, B. Lundberg⁷⁸,
 J. Lundberg^{145a,145b}, O. Lundberg^{145a,145b}, J. Lundquist³⁵, M. Lungwitz⁸⁰, D. Lynn²⁴, E. Lytken⁷⁸,
 H. Ma²⁴, L.L. Ma¹⁷², G. Maccarrone⁴⁶, A. Macchiolo⁹⁸, B. Maček⁷³, J. Machado Miguens^{123a},
 R. Mackeprang³⁵, R.J. Madaras¹⁴, H.J. Maddocks⁷⁰, W.F. Mader⁴³, R. Maenner^{57c}, T. Maeno²⁴,
 P. Mättig¹⁷⁴, S. Mättig⁴¹, L. Magnoni²⁹, E. Magradze⁵³, K. Mahboubi⁴⁷, S. Mahmoud⁷², G. Mahout¹⁷,
 C. Maiani¹³⁵, C. Maidantchik^{23a}, A. Maio^{123a,b}, S. Majewski²⁴, Y. Makida⁶⁴, N. Makovec¹¹⁴, P. Mal¹³⁵,
 B. Malaescu²⁹, Pa. Malecki³⁸, P. Malecki³⁸, V.P. Maleev¹²⁰, F. Malek⁵⁴, U. Mallik⁶¹, D. Malon⁵,
 C. Malone¹⁴², S. Maltezos⁹, V. Malyshev¹⁰⁶, S. Malyukov²⁹, R. Mameghani⁹⁷, J. Mamuzic^{12b},
 A. Manabe⁶⁴, L. Mandelli^{88a}, I. Mandić⁷³, R. Mandrysch¹⁵, J. Maneira^{123a}, P.S. Mangedard⁸⁷,
 L. Manhaes de Andrade Filho^{23b}, J.A. Manjarres Ramos¹³⁵, A. Mann⁵³, P.M. Manning¹³⁶,
 A. Manousakis-Katsikakis⁸, B. Mansoulie¹³⁵, A. Mapelli²⁹, L. Mapelli²⁹, L. March⁷⁹, J.F. Marchand²⁸,
 F. Marchese^{132a,132b}, G. Marchiori⁷⁷, M. Marcisovsky¹²⁴, C.P. Marino¹⁶⁸, F. Marroquim^{23a}, Z. Marshall²⁹,
 F.K. Martens¹⁵⁷, L.F. Marti¹⁶, S. Marti-Garcia¹⁶⁶, B. Martin²⁹, B. Martin⁸⁷, J.P. Martin⁹², T.A. Martin¹⁷,
 V.J. Martin⁴⁵, B. Martin dit Latour⁴⁸, S. Martin-Haugh¹⁴⁸, M. Martinez¹¹, V. Martinez Outschoorn⁵⁶,
 A.C. Martyniuk¹⁶⁸, M. Marx⁸¹, F. Marzano^{131a}, A. Marzin¹¹⁰, L. Masetti⁸⁰, T. Mashimo¹⁵⁴,
 R. Mashinistov⁹³, J. Masik⁸¹, A.L. Maslennikov¹⁰⁶, I. Massa^{19a,19b}, G. Massaro¹⁰⁴, N. Massol⁴,
 P. Mastrandrea¹⁴⁷, A. Mastroberardino^{36a,36b}, T. Masubuchi¹⁵⁴, P. Matricon¹¹⁴, H. Matsunaga¹⁵⁴,
 T. Matsushita⁶⁵, C. Mattravers^{117,c}, J. Maurer⁸², S.J. Maxfield⁷², A. Mayne¹³⁸, R. Mazini¹⁵⁰, M. Mazur²⁰,
 L. Mazzaferro^{132a,132b}, M. Mazzanti^{88a}, S.P. Mc Kee⁸⁶, A. McCarn¹⁶⁴, R.L. McCarthy¹⁴⁷, T.G. McCarthy²⁸,
 N.A. McCubbin¹²⁸, K.W. McFarlane^{55,*}, J.A. MCFayden¹³⁸, G. Mchedlidze^{50b}, T. McLaughlan¹⁷,

S.J. McMahon¹²⁸, R.A. McPherson^{168,k}, A. Meade⁸³, J. Mechnich¹⁰⁴, M. Mechtel¹⁷⁴, M. Medinnis⁴¹,
 R. Meera-Lebbai¹¹⁰, T. Meguro¹¹⁵, R. Mehdiyev⁹², S. Mehlhase³⁵, A. Mehta⁷², K. Meier^{57a},
 B. Meirose⁷⁸, C. Melachrinou³⁰, B.R. Mellado Garcia¹⁷², F. Meloni^{88a,88b}, L. Mendoza Navas¹⁶¹,
 Z. Meng^{150,u}, A. Mengarelli^{19a,19b}, S. Menke⁹⁸, E. Meoni¹⁶⁰, K.M. Mercurio⁵⁶, P. Mermod⁴⁸,
 L. Merola^{101a,101b}, C. Meroni^{88a}, F.S. Merritt³⁰, H. Merritt¹⁰⁸, A. Messina^{29,y}, J. Metcalfe¹⁰²,
 A.S. Mete¹⁶², C. Meyer⁸⁰, C. Meyer³⁰, J.-P. Meyer¹³⁵, J. Meyer¹⁷³, J. Meyer⁵³, T.C. Meyer²⁹,
 W.T. Meyer⁶², J. Miao^{32d}, S. Michal²⁹, L. Micu^{25a}, R.P. Middleton¹²⁸, S. Migas⁷², L. Mijović¹³⁵,
 G. Mikenberg¹⁷¹, M. Mikestikova¹²⁴, M. Mikuž⁷³, D.W. Miller³⁰, R.J. Miller⁸⁷, W.J. Mills¹⁶⁷, C. Mills⁵⁶,
 A. Milov¹⁷¹, D.A. Milstead^{145a,145b}, D. Milstein¹⁷¹, A.A. Minaenko¹²⁷, M. Miñano Moya¹⁶⁶,
 I.A. Minashvili⁶³, A.I. Mincer¹⁰⁷, B. Mindur³⁷, M. Mineev⁶³, Y. Ming¹⁷², L.M. Mir¹¹, G. Mirabelli^{131a},
 J. Mitrevski¹³⁶, V.A. Mitsou¹⁶⁶, S. Mitsui⁶⁴, P.S. Miyagawa¹³⁸, J.U. Mjörnmark⁷⁸, T. Moa^{145a,145b},
 V. Moeller²⁷, K. Mönig⁴¹, N. Möser²⁰, S. Mohapatra¹⁴⁷, W. Mohr⁴⁷, R. Moles-Valls¹⁶⁶, J. Monk⁷⁶,
 E. Monnier⁸², J. Montejo Berlingen¹¹, F. Monticelli⁶⁹, S. Monzani^{19a,19b}, R.W. Moore², G.F. Moorhead⁸⁵,
 C. Mora Herrera⁴⁸, A. Moraes⁵², N. Morange¹³⁵, J. Morel⁵³, G. Morello^{36a,36b}, D. Moreno⁸⁰,
 M. Moreno Llácer¹⁶⁶, P. Morettini^{49a}, M. Morgenstern⁴³, M. Morii⁵⁶, A.K. Morley²⁹, G. Mornacchi²⁹,
 J.D. Morris⁷⁴, L. Morvaj¹⁰⁰, H.G. Moser⁹⁸, M. Mosidze^{50b}, J. Moss¹⁰⁸, R. Mount¹⁴², E. Mountricha^{9,z},
 S.V. Mouraviev^{93,*}, E.J.W. Moyse⁸³, F. Mueller^{57a}, J. Mueller¹²², K. Mueller²⁰, T.A. Müller⁹⁷,
 T. Mueller⁸⁰, D. Muenstermann²⁹, Y. Munwes¹⁵², W.J. Murray¹²⁸, I. Mussche¹⁰⁴, E. Musto^{101a,101b},
 A.G. Myagkov¹²⁷, M. Myska¹²⁴, J. Nadal¹¹, K. Nagai¹⁵⁹, K. Nagano⁶⁴, A. Nagarkar¹⁰⁸, Y. Nagasaka⁵⁸,
 M. Nagel⁹⁸, A.M. Nairz²⁹, Y. Nakahama²⁹, K. Nakamura¹⁵⁴, T. Nakamura¹⁵⁴, I. Nakano¹⁰⁹, G. Nanava²⁰,
 A. Napier¹⁶⁰, R. Narayan^{57b}, M. Nash^{76,c}, T. Nattermann²⁰, T. Naumann⁴¹, G. Navarro¹⁶¹, H.A. Neal⁸⁶,
 P.Yu. Nechaeva⁹³, T.J. Neep⁸¹, A. Negri^{118a,118b}, G. Negri²⁹, M. Negrini^{19a}, S. Nektarijevic⁴⁸,
 A. Nelson¹⁶², T.K. Nelson¹⁴², S. Nemecek¹²⁴, P. Nemethy¹⁰⁷, A.A. Nepomuceno^{23a}, M. Nessi^{29,aa},
 M.S. Neubauer¹⁶⁴, A. Neusiedl⁸⁰, R.M. Neves¹⁰⁷, P. Nevski²⁴, P.R. Newman¹⁷, V. Nguyen Thi Hong¹³⁵,
 R.B. Nickerson¹¹⁷, R. Nicolaidou¹³⁵, B. Nicquevert²⁹, F. Niedercorn¹¹⁴, J. Nielsen¹³⁶, N. Nikiforou³⁴,
 A. Nikiforov¹⁵, V. Nikolaenko¹²⁷, I. Nikolic-Audit⁷⁷, K. Nikolics⁴⁸, K. Nikolopoulos¹⁷, H. Nilsen⁴⁷,
 P. Nilsson⁷, Y. Ninomiya¹⁵⁴, A. Nisati^{131a}, R. Nisius⁹⁸, T. Nobe¹⁵⁶, L. Nodulman⁵, M. Nomachi¹¹⁵,
 I. Nomidis¹⁵³, S. Norberg¹¹⁰, M. Nordberg²⁹, P.R. Norton¹²⁸, J. Novakova¹²⁵, M. Nozaki⁶⁴, L. Nozka¹¹²,
 I.M. Nugent^{158a}, A.-E. Nuncio-Quiroz²⁰, G. Nunes Hanninger⁸⁵, T. Nunnemann⁹⁷, E. Nurse⁷⁶,
 B.J. O'Brien⁴⁵, S.W. O'Neale^{17,*}, D.C. O'Neil¹⁴¹, V. O'Shea⁵², L.B. Oakes⁹⁷, F.G. Oakham^{28,d},
 H. Oberlack⁹⁸, J. Ocariz⁷⁷, A. Ochi⁶⁵, S. Oda⁶⁸, S. Odaka⁶⁴, J. Odier⁸², H. Ogren⁵⁹, A. Oh⁸¹, S.H. Oh⁴⁴,
 C.C. Ohm²⁹, T. Ohshima¹⁰⁰, H. Okawa²⁴, Y. Okumura³⁰, T. Okuyama¹⁵⁴, A. Olariu^{25a}, A.G. Olchevski⁶³,
 S.A. Olivares Pino^{31a}, M. Oliveira^{123a,h}, D. Oliveira Damazio²⁴, E. Oliver Garcia¹⁶⁶, D. Olivito¹¹⁹,
 A. Olszewski³⁸, J. Olszowska³⁸, A. Onofre^{123a,ab}, P.U.E. Onyisi³⁰, C.J. Oram^{158a}, M.J. Oreglia³⁰,
 Y. Oren¹⁵², D. Orestano^{133a,133b}, N. Orlando^{71a,71b}, I. Orlov¹⁰⁶, C. Oropeza Barrera⁵², R.S. Orr¹⁵⁷,
 B. Osculati^{49a,49b}, R. Ospanov¹¹⁹, C. Osuna¹¹, G. Otero y Garzon²⁶, J.P. Ottersbach¹⁰⁴, M. Ouchrif^{134d},
 E.A. Ouellette¹⁶⁸, F. Ould-Saada¹¹⁶, A. Ouraou¹³⁵, Q. Ouyang^{32a}, A. Ovcharova¹⁴, M. Owen⁸¹,
 S. Owen¹³⁸, V.E. Ozcan^{18a}, N. Ozturk⁷, A. Pacheco Pages¹¹, C. Padilla Aranda¹¹, S. Pagan Griso¹⁴,
 E. Paganis¹³⁸, C. Pahl⁹⁸, F. Paige²⁴, P. Pais⁸³, K. Pajchel¹¹⁶, G. Palacino^{158b}, C.P. Paleari⁶, S. Palestini²⁹,
 D. Pallin³³, A. Palma^{123a}, J.D. Palmer¹⁷, Y.B. Pan¹⁷², E. Panagiotopoulou⁹, P. Pani¹⁰⁴, N. Panikashvili⁸⁶,
 S. Panitkin²⁴, D. Pantea^{25a}, A. Papadelis^{145a}, Th.D. Papadopoulou⁹, A. Paramonov⁵,
 D. Paredes Hernandez³³, W. Park^{24,ac}, M.A. Parker²⁷, F. Parodi^{49a,49b}, J.A. Parsons³⁴, U. Parzefall⁴⁷,
 S. Pashapour⁵³, E. Pasqualucci^{131a}, S. Passaggio^{49a}, A. Passeri^{133a}, F. Pastore^{133a,133b,*}, Fr. Pastore⁷⁵,
 G. Pásztor^{48,ad}, S. Pataraja¹⁷⁴, N. Patel¹⁴⁹, J.R. Pater⁸¹, S. Patricelli^{101a,101b}, T. Pauly²⁹, M. Pecsny^{143a},
 M.I. Pedraza Morales¹⁷², S.V. Peleganchuk¹⁰⁶, D. Pelikan¹⁶⁵, H. Peng^{32b}, B. Penning³⁰, A. Penson³⁴,
 J. Penwell⁵⁹, M. Perantoni^{23a}, K. Perez^{34,ae}, T. Perez Cavalcanti⁴¹, E. Perez Codina^{158a},
 M.T. Pérez García-Estañ¹⁶⁶, V. Perez Reale³⁴, L. Perini^{88a,88b}, H. Pernegger²⁹, R. Perrino^{71a}, P. Perrodo⁴,
 V.D. Peshekhonov⁶³, K. Peters²⁹, B.A. Petersen²⁹, J. Petersen²⁹, T.C. Petersen³⁵, E. Petit⁴, A. Petridis¹⁵³,
 C. Petridou¹⁵³, E. Petrolu^{131a}, F. Petrucci^{133a,133b}, D. Petschull⁴¹, M. Petteni¹⁴¹, R. Pezoa^{31b}, A. Phan⁸⁵,
 P.W. Phillips¹²⁸, G. Piacquadio²⁹, A. Picazio⁴⁸, E. Piccaro⁷⁴, M. Piccinini^{19a,19b}, S.M. Piec⁴¹, R. Piegai²⁶,
 D.T. Pignotti¹⁰⁸, J.E. Pilcher³⁰, A.D. Pilkington⁸¹, J. Pina^{123a,b}, M. Pinamonti^{163a,163c}, A. Pinder¹¹⁷,
 J.L. Pinfold², B. Pinto^{123a}, C. Pizio^{88a,88b}, M. Plamondon¹⁶⁸, M.-A. Pleier²⁴, E. Plotnikova⁶³,

A. Poblaguev²⁴, S. Poddar^{57a}, F. Podlyski³³, L. Poggioli¹¹⁴, M. Pohl⁴⁸, G. Polesello^{118a},
 A. Policicchio^{36a,36b}, A. Polini^{19a}, J. Poll⁷⁴, V. Polychronakos²⁴, D. Pomeroy²², K. Pommès²⁹,
 L. Pontecorvo^{131a}, B.G. Pope⁸⁷, G.A. Popeneciu^{25a}, D.S. Popovic^{12a}, A. Poppleton²⁹, X. Portell Bueso²⁹,
 G.E. Pospelov⁹⁸, S. Pospisil¹²⁶, I.N. Potrap⁹⁸, C.J. Potter¹⁴⁸, C.T. Potter¹¹³, G. Poulard²⁹, J. Poveda⁵⁹,
 V. Pozdnyakov⁶³, R. Prabhu⁷⁶, P. Pralavorio⁸², A. Pranko¹⁴, S. Prasad²⁹, R. Pravahan²⁴, S. Prell⁶²,
 K. Pretzl¹⁶, D. Price⁵⁹, J. Price⁷², L.E. Price⁵, D. Prieur¹²², M. Primavera^{71a}, K. Prokofiev¹⁰⁷,
 F. Prokoshin^{31b}, S. Protopopescu²⁴, J. Proudfoot⁵, X. Prudent⁴³, M. Przybycien³⁷, H. Przysiezniak⁴,
 S. Psoroulas²⁰, E. Ptacek¹¹³, E. Pueschel⁸³, J. Purdham⁸⁶, M. Purohit^{24,ac}, P. Puzo¹¹⁴, Y. Pylypchenko⁶¹,
 J. Qian⁸⁶, A. Quadt⁵³, D.R. Quarrie¹⁴, W.B. Quayle¹⁷², F. Quinonez^{31a}, M. Raas¹⁰³, V. Radescu⁴¹,
 P. Radloff¹¹³, T. Rador^{18a}, F. Ragusa^{88a,88b}, G. Rahal¹⁷⁷, A.M. Rahimi¹⁰⁸, D. Rahm²⁴, S. Rajagopalan²⁴,
 M. Rammensee⁴⁷, M. Rammes¹⁴⁰, A.S. Randle-Conde³⁹, K. Randrianarivony²⁸, F. Rauscher⁹⁷,
 T.C. Rave⁴⁷, M. Raymond²⁹, A.L. Read¹¹⁶, D.M. Rebuffi^{118a,118b}, A. Redelbach¹⁷³, G. Redlinger²⁴,
 R. Reece¹¹⁹, K. Reeves⁴⁰, E. Reinherz-Aronis¹⁵², A. Reinsch¹¹³, I. Reisinger⁴², C. Rembser²⁹, Z.L. Ren¹⁵⁰,
 A. Renaud¹¹⁴, M. Rescigno^{131a}, S. Resconi^{88a}, B. Resende¹³⁵, P. Reznicek⁹⁷, R. Rezvani¹⁵⁷, R. Richter⁹⁸,
 E. Richter-Was^{4,af}, M. Ridel⁷⁷, M. Rijpstra¹⁰⁴, M. Rijssenbeek¹⁴⁷, A. Rimoldi^{118a,118b}, L. Rinaldi^{19a},
 R.R. Rios³⁹, I. Riu¹¹, G. Rivoltella^{88a,88b}, F. Rizatdinova¹¹¹, E. Rizvi⁷⁴, S.H. Robertson^{84,k},
 A. Robichaud-Veronneau¹¹⁷, D. Robinson²⁷, J.E.M. Robinson⁸¹, A. Robson⁵², J.G. Rocha de Lima¹⁰⁵,
 C. Roda^{121a,121b}, D. Roda Dos Santos²⁹, A. Roe⁵³, S. Roe²⁹, O. Røhne¹¹⁶, S. Rolli¹⁶⁰, A. Romaniouk⁹⁵,
 M. Romano^{19a,19b}, G. Romeo²⁶, E. Romero Adam¹⁶⁶, L. Roos⁷⁷, E. Ros¹⁶⁶, S. Rosati^{131a}, K. Rosbach⁴⁸,
 A. Rose¹⁴⁸, M. Rose⁷⁵, G.A. Rosenbaum¹⁵⁷, E.I. Rosenberg⁶², P.L. Rosendahl¹³, O. Rosenthal¹⁴⁰,
 L. Rossetlet⁴⁸, V. Rossetti¹¹, E. Rossi^{131a,131b}, L.P. Rossi^{49a}, M. Rotaru^{25a}, I. Roth¹⁷¹, J. Rothberg¹³⁷,
 D. Rousseau¹¹⁴, C.R. Royon¹³⁵, A. Rozanov⁸², Y. Rozen¹⁵¹, X. Ruan^{32a,ag}, F. Rubbo¹¹, I. Rubinskiy⁴¹,
 B. Ruckert⁹⁷, N. Ruckstuhl¹⁰⁴, V.I. Rud⁹⁶, C. Rudolph⁴³, G. Rudolph⁶⁰, F. Rühr⁶, A. Ruiz-Martinez⁶²,
 L. Rummyantsev⁶³, Z. Rurikova⁴⁷, N.A. Rusakovich⁶³, J.P. Rutherford⁶, C. Ruwiedel^{14,*}, P. Ruzicka¹²⁴,
 Y.F. Ryabov¹²⁰, P. Ryan⁸⁷, M. Rybar¹²⁵, G. Rybkin¹¹⁴, N.C. Ryder¹¹⁷, A.F. Saavedra¹⁴⁹, I. Sadeh¹⁵²,
 H.F.-W. Sadrozinski¹³⁶, R. Sadykov⁶³, F. Safai Tehrani^{131a}, H. Sakamoto¹⁵⁴, G. Salamanna⁷⁴,
 A. Salamon^{132a}, M. Saleem¹¹⁰, D. Salek²⁹, D. Salihagic⁹⁸, A. Salnikov¹⁴², J. Salt¹⁶⁶,
 B.M. Salvachua Ferrando⁵, D. Salvatore^{36a,36b}, F. Salvatore¹⁴⁸, A. Salvucci¹⁰³, A. Salzburger²⁹,
 D. Sampsonidis¹⁵³, B.H. Samset¹¹⁶, A. Sanchez^{101a,101b}, V. Sanchez Martinez¹⁶⁶, H. Sandaker¹³,
 H.G. Sander⁸⁰, M.P. Sanders⁹⁷, M. Sandhoff¹⁷⁴, T. Sandoval²⁷, C. Sandoval¹⁶¹, R. Sandstroem⁹⁸,
 D.P.C. Sankey¹²⁸, A. Sansoni⁴⁶, C. Santamarina Rios⁸⁴, C. Santoni³³, R. Santonico^{132a,132b}, H. Santos^{123a},
 J.G. Saraiva^{123a}, T. Sarangi¹⁷², E. Sarkisyan-Grinbaum⁷, F. Sarri^{121a,121b}, G. Sartisohn¹⁷⁴, O. Sasaki⁶⁴,
 N. Sasao⁶⁶, I. Satsounkevitch⁸⁹, G. Sauvage^{4,*}, E. Sauvan⁴, J.B. Sauvan¹¹⁴, P. Savard^{157,d}, V. Savinov¹²²,
 D.O. Savu²⁹, L. Sawyer^{24,m}, D.H. Saxon⁵², J. Saxon¹¹⁹, C. Sbarra^{19a}, A. Sbrizzi^{19a,19b},
 D.A. Scannicchio¹⁶², M. Scarcella¹⁴⁹, J. Schaarschmidt¹¹⁴, P. Schacht⁹⁸, D. Schaefer¹¹⁹, U. Schäfer⁸⁰,
 S. Schaepe²⁰, S. Schaezel^{57b}, A.C. Schaffer¹¹⁴, D. Schaile⁹⁷, R.D. Schamberger¹⁴⁷, A.G. Schamov¹⁰⁶,
 V. Scharf^{57a}, V.A. Schegelsky¹²⁰, D. Scheirich⁸⁶, M. Schernau¹⁶², M.I. Scherzer³⁴, C. Schiavi^{49a,49b},
 J. Schieck⁹⁷, M. Schioppa^{36a,36b}, S. Schlenker²⁹, E. Schmidt⁴⁷, K. Schmieden²⁰, C. Schmitt⁸⁰,
 S. Schmitt^{57b}, M. Schmitz²⁰, B. Schneider¹⁶, U. Schnoor⁴³, A. Schoening^{57b}, A.L.S. Schorlemmer⁵³,
 M. Schott²⁹, D. Schouten^{158a}, J. Schovancova¹²⁴, M. Schram⁸⁴, C. Schroeder⁸⁰, N. Schroer^{57c},
 M.J. Schultens²⁰, J. Schultes¹⁷⁴, H.-C. Schultz-Coulon^{57a}, H. Schulz¹⁵, M. Schumacher⁴⁷,
 B.A. Schumm¹³⁶, Ph. Schune¹³⁵, C. Schwanenberger⁸¹, A. Schwartzman¹⁴², Ph. Schwemling⁷⁷,
 R. Schwienhorst⁸⁷, R. Schwierz⁴³, J. Schwindling¹³⁵, T. Schwindt²⁰, M. Schwoerer⁴, G. Sciolla²²,
 W.G. Scott¹²⁸, J. Searcy¹¹³, G. Sedov⁴¹, E. Sedykh¹²⁰, S.C. Seidel¹⁰², A. Seiden¹³⁶, F. Seifert⁴³,
 J.M. Seixas^{23a}, G. Sekhniaidze^{101a}, S.J. Sekula³⁹, K.E. Selbach⁴⁵, D.M. Seliverstov¹²⁰, B. Sellden^{145a},
 G. Sellers⁷², M. Seman^{143b}, N. Semprini-Cesari^{19a,19b}, C. Serfon⁹⁷, L. Serin¹¹⁴, L. Serkin⁵³, R. Seuster⁹⁸,
 H. Severini¹¹⁰, A. Sfyrla²⁹, E. Shabalina⁵³, M. Shamim¹¹³, L.Y. Shan^{32a}, J.T. Shank²¹, Q.T. Shao⁸⁵,
 M. Shapiro¹⁴, P.B. Shatalov⁹⁴, K. Shaw^{163a,163c}, D. Sherman¹⁷⁵, P. Sherwood⁷⁶, A. Shibata¹⁰⁷,
 S. Shimizu²⁹, M. Shimojima⁹⁹, T. Shin⁵⁵, M. Shiyakova⁶³, A. Shmeleva⁹³, M.J. Shochet³⁰, D. Short¹¹⁷,
 S. Shrestha⁶², E. Shulga⁹⁵, M.A. Shupe⁶, P. Sicho¹²⁴, A. Sidoti^{131a}, F. Siegert⁴⁷, Dj. Sijacki^{12a},
 O. Silbert¹⁷¹, J. Silva^{123a}, Y. Silver¹⁵², D. Silverstein¹⁴², S.B. Silverstein^{145a}, V. Simak¹²⁶, O. Simard¹³⁵,
 Lj. Simic^{12a}, S. Simion¹¹⁴, E. Simioni⁸⁰, B. Simmons⁷⁶, R. Simoniello^{88a,88b}, M. Simonyan³⁵,

P. Sinervo¹⁵⁷, N.B. Sinev¹¹³, V. Sipica¹⁴⁰, G. Siragusa¹⁷³, A. Sircar²⁴, A.N. Sisakyan^{63,*},
 S.Yu. Sivoklokov⁹⁶, J. Sjölin^{145a,145b}, T.B. Sjursen¹³, L.A. Skinnari¹⁴, H.P. Skottowe⁵⁶, K. Skovpen¹⁰⁶,
 P. Skubic¹¹⁰, M. Slater¹⁷, T. Slavicek¹²⁶, K. Sliwa¹⁶⁰, V. Smakhtin¹⁷¹, B.H. Smart⁴⁵, S.Yu. Smirnov⁹⁵,
 Y. Smirnov⁹⁵, L.N. Smirnova⁹⁶, O. Smirnova⁷⁸, B.C. Smith⁵⁶, D. Smith¹⁴², K.M. Smith⁵²,
 M. Smizanska⁷⁰, K. Smolek¹²⁶, A.A. Snesarev⁹³, S.W. Snow⁸¹, J. Snow¹¹⁰, S. Snyder²⁴, R. Sobie^{168,k},
 J. Sodomka¹²⁶, A. Soffer¹⁵², C.A. Solans¹⁶⁶, M. Solar¹²⁶, J. Solc¹²⁶, E.Yu. Soldatov⁹⁵, U. Soldevila¹⁶⁶,
 E. Solfaroli Camillocci^{131a,131b}, A.A. Solodkov¹²⁷, O.V. Solovyanov¹²⁷, N. Soni⁸⁵, V. Sopko¹²⁶,
 B. Sopko¹²⁶, M. Sosebee⁷, R. Soualah^{163a,163c}, A. Soukharev¹⁰⁶, S. Spagnolo^{71a,71b}, F. Spanò⁷⁵,
 R. Spighi^{19a}, G. Spigo²⁹, R. Spiwoks²⁹, M. Spousta^{125,ah}, T. Spreitzer¹⁵⁷, B. Spurlock⁷, R.D. St. Denis⁵²,
 J. Stahlman¹¹⁹, R. Stamen^{57a}, E. Stanecka³⁸, R.W. Stanek⁵, C. Stanescu^{133a}, M. Stanescu-Bellu⁴¹,
 S. Stapnes¹¹⁶, E.A. Starchenko¹²⁷, J. Stark⁵⁴, P. Staroba¹²⁴, P. Starovoitov⁴¹, R. Staszewski³⁸,
 A. Staude⁹⁷, P. Stavina^{143a,*}, G. Steele⁵², P. Steinbach⁴³, P. Steinberg²⁴, I. Stekl¹²⁶, B. Stelzer¹⁴¹,
 H.J. Stelzer⁸⁷, O. Stelzer-Chilton^{158a}, H. Stenzel⁵¹, S. Stern⁹⁸, G.A. Stewart²⁹, J.A. Stillings²⁰,
 M.C. Stockton⁸⁴, K. Stoerig⁴⁷, G. Stoicea^{25a}, S. Stonjek⁹⁸, P. Strachota¹²⁵, A.R. Stradling⁷,
 A. Straessner⁴³, J. Strandberg¹⁴⁶, S. Strandberg^{145a,145b}, A. Strandlie¹¹⁶, M. Strang¹⁰⁸, E. Strauss¹⁴²,
 M. Strauss¹¹⁰, P. Strizenec^{143b}, R. Ströhmer¹⁷³, D.M. Strom¹¹³, J.A. Strong^{75,*}, R. Stroynowski³⁹,
 J. Strube¹²⁸, B. Stugu¹³, I. Stumer^{24,*}, J. Stupak¹⁴⁷, P. Sturm¹⁷⁴, N.A. Styles⁴¹, D.A. Soh^{150,w}, D. Su¹⁴²,
 HS. Subramania², A. Succurro¹¹, Y. Sugaya¹¹⁵, C. Suhr¹⁰⁵, M. Suk¹²⁵, V.V. Sulin⁹³, S. Sultansoy^{3d},
 T. Sumida⁶⁶, X. Sun⁵⁴, J.E. Sundermann⁴⁷, K. Suruliz¹³⁸, G. Susinno^{36a,36b}, M.R. Sutton¹⁴⁸, Y. Suzuki⁶⁴,
 Y. Suzuki⁶⁵, M. Svatos¹²⁴, S. Swedish¹⁶⁷, I. Sykora^{143a}, T. Sykora¹²⁵, J. Sánchez¹⁶⁶, D. Ta¹⁰⁴,
 K. Tackmann⁴¹, A. Taffard¹⁶², R. Tafirout^{158a}, N. Taiblum¹⁵², Y. Takahashi¹⁰⁰, H. Takai²⁴,
 R. Takashima⁶⁷, H. Takeda⁶⁵, T. Takeshita¹³⁹, Y. Takubo⁶⁴, M. Talby⁸², A. Talyshv^{106,f}, M.C. Tamsett²⁴,
 J. Tanaka¹⁵⁴, R. Tanaka¹¹⁴, S. Tanaka¹³⁰, S. Tanaka⁶⁴, A.J. Tanasijczuk¹⁴¹, K. Tani⁶⁵, N. Tannoury⁸²,
 S. Tapprogge⁸⁰, D. Tardif¹⁵⁷, S. Tarem¹⁵¹, F. Tarrade²⁸, G.F. Tartarelli^{88a}, P. Tas¹²⁵, M. Tasevsky¹²⁴,
 E. Tassi^{36a,36b}, M. Tatarkhanov¹⁴, Y. Tayalati^{134d}, C. Taylor⁷⁶, F.E. Taylor⁹¹, G.N. Taylor⁸⁵, W. Taylor^{158b},
 M. Teinturier¹¹⁴, M. Teixeira Dias Castanheira⁷⁴, P. Teixeira-Dias⁷⁵, K.K. Temming⁴⁷, H. Ten Kate²⁹,
 P.K. Teng¹⁵⁰, S. Terada⁶⁴, K. Terashi¹⁵⁴, J. Terron⁷⁹, M. Testa⁴⁶, R.J. Teuscher^{157,k}, J. Therhaag²⁰,
 T. Theveneaux-Pelzer⁷⁷, S. Thoma⁴⁷, J.P. Thomas¹⁷, E.N. Thompson³⁴, P.D. Thompson¹⁷,
 P.D. Thompson¹⁵⁷, A.S. Thompson⁵², L.A. Thomsen³⁵, E. Thomson¹¹⁹, M. Thomson²⁷, W.M. Thong⁸⁵,
 R.P. Thun⁸⁶, F. Tian³⁴, M.J. Tibbetts¹⁴, T. Tic¹²⁴, V.O. Tikhomirov⁹³, Y.A. Tikhonov^{106,f}, S. Timoshenko⁹⁵,
 P. Tipton¹⁷⁵, S. Tisserant⁸², T. Todorov⁴, S. Todorova-Nova¹⁶⁰, B. Toggerson¹⁶², J. Tojo⁶⁸, S. Tokár^{143a},
 K. Tokushuku⁶⁴, K. Tollefson⁸⁷, M. Tomoto¹⁰⁰, L. Tompkins³⁰, K. Toms¹⁰², A. Tonoyan¹³, C. Topfel¹⁶,
 N.D. Topilin⁶³, I. Torchiani²⁹, E. Torrence¹¹³, H. Torres⁷⁷, E. Torró Pastor¹⁶⁶, J. Toth^{82,ad}, F. Touchard⁸²,
 D.R. Tovey¹³⁸, T. Trefzger¹⁷³, L. Tremblet²⁹, A. Tricoli²⁹, I.M. Trigger^{158a}, S. Trincaz-Duvold⁷⁷,
 M.F. Tripiana⁶⁹, N. Triplett²⁴, W. Trischuk¹⁵⁷, B. Trocmé⁵⁴, C. Troncon^{88a}, M. Trottier-McDonald¹⁴¹,
 M. Trzebinski³⁸, A. Trzupek³⁸, C. Tsarouchas²⁹, J.C.-L. Tseng¹¹⁷, M. Tsiakiris¹⁰⁴, P.V. Tsiarehka⁸⁹,
 D. Tsiou^{4,ai}, G. Tsipolitis⁹, S. Tsiskaridze¹¹, V. Tsiskaridze⁴⁷, E.G. Tskhadadze^{50a}, I.I. Tsukerman⁹⁴,
 V. Tsulaia¹⁴, J.-W. Tsung²⁰, S. Tsuno⁶⁴, D. Tsybychev¹⁴⁷, A. Tua¹³⁸, A. Tudorache^{25a}, V. Tudorache^{25a},
 J.M. Tuggle³⁰, M. Turala³⁸, D. Turecek¹²⁶, I. Turk Cakir^{3e}, E. Turlay¹⁰⁴, R. Turra^{88a,88b}, P.M. Tuts³⁴,
 A. Tykhonov⁷³, M. Tylmad^{145a,145b}, M. Tyndel¹²⁸, G. Tzanakos⁸, K. Uchida²⁰, I. Ueda¹⁵⁴, R. Ueno²⁸,
 M. Uglan¹³, M. Uhlenbrock²⁰, M. Uhrmacher⁵³, F. Ukegawa¹⁵⁹, G. Unal²⁹, A. Undrus²⁴, G. Unel¹⁶²,
 Y. Unno⁶⁴, D. Urbaniec³⁴, G. Usai⁷, M. Uslenghi^{118a,118b}, L. Vacavant⁸², V. Vacek¹²⁶, B. Vachon⁸⁴,
 S. Vahsen¹⁴, J. Valenta¹²⁴, S. Valentinetti^{19a,19b}, A. Valero¹⁶⁶, S. Valkar¹²⁵, E. Valladolid Gallego¹⁶⁶,
 S. Vallecorsa¹⁵¹, J.A. Valls Ferrer¹⁶⁶, P.C. Van Der Deijl¹⁰⁴, R. van der Geer¹⁰⁴, H. van der Graaf¹⁰⁴,
 E. van der Kraaij¹⁰⁴, R. Van Der Leeuw¹⁰⁴, E. van der Poel¹⁰⁴, D. van der Ster²⁹, N. van Eldik²⁹,
 P. van Gemmeren⁵, I. van Vulpen¹⁰⁴, M. Vanadia⁹⁸, W. Vandelli²⁹, A. Vaniachine⁵, P. Vankov⁴¹,
 F. Vannucci⁷⁷, R. Vari^{131a}, T. Varol⁸³, D. Varouchas¹⁴, A. Vartapetian⁷, K.E. Varvell¹⁴⁹,
 V.I. Vassilakopoulos⁵⁵, F. Vazeille³³, T. Vazquez Schroeder⁵³, G. Vegni^{88a,88b}, J.J. Veillet¹¹⁴, F. Veloso^{123a},
 R. Veness²⁹, S. Veneziano^{131a}, A. Ventura^{71a,71b}, D. Ventura⁸³, M. Venturi⁴⁷, N. Venturi¹⁵⁷,
 V. Vercesi^{118a}, M. Verducci¹³⁷, W. Verkerke¹⁰⁴, J.C. Vermeulen¹⁰⁴, A. Vest⁴³, M.C. Vetterli^{141,d},
 I. Vichou¹⁶⁴, T. Vickey^{144b,aj}, O.E. Vickey Boeriu^{144b}, G.H.A. Viehhauser¹¹⁷, S. Viel¹⁶⁷, M. Villa^{19a,19b},
 M. Villaplana Perez¹⁶⁶, E. Vilucchi⁴⁶, M.G. Vincter²⁸, E. Vinek²⁹, V.B. Vinogradov⁶³, M. Virchaux^{135,*},

J. Virzi¹⁴, O. Vitells¹⁷¹, M. Viti⁴¹, I. Vivarelli⁴⁷, F. Vives Vaque², S. Vlachos⁹, D. Vladoiu⁹⁷, M. Vlasak¹²⁶, A. Vogel²⁰, P. Vokac¹²⁶, G. Volpi⁴⁶, M. Volpi⁸⁵, G. Volpini^{88a}, H. von der Schmitt⁹⁸, J. von Loeben⁹⁸, H. von Radziewski⁴⁷, E. von Toerne²⁰, V. Vorobel¹²⁵, V. Vorwerk¹¹, M. Vos¹⁶⁶, R. Voss²⁹, T.T. Voss¹⁷⁴, J.H. Vossebeld⁷², N. Vranjes¹³⁵, M. Vranjes Milosavljevic¹⁰⁴, V. Vrba¹²⁴, M. Vreeswijk¹⁰⁴, T. Vu Anh⁴⁷, R. Vuillemet²⁹, I. Vukotic³⁰, W. Wagner¹⁷⁴, P. Wagner¹¹⁹, H. Wahlen¹⁷⁴, S. Wahrmund⁴³, J. Wakabayashi¹⁰⁰, S. Walch⁸⁶, J. Walder⁷⁰, R. Walker⁹⁷, W. Walkowiak¹⁴⁰, R. Wall¹⁷⁵, P. Waller⁷², B. Walsh¹⁷⁵, C. Wang⁴⁴, H. Wang¹⁷², H. Wang^{32b,ak}, J. Wang¹⁵⁰, J. Wang⁵⁴, R. Wang¹⁰², S.M. Wang¹⁵⁰, T. Wang²⁰, A. Warburton⁸⁴, C.P. Ward²⁷, M. Warsinsky⁴⁷, A. Washbrook⁴⁵, C. Wasicki⁴¹, I. Watanabe⁶⁵, P.M. Watkins¹⁷, A.T. Watson¹⁷, I.J. Watson¹⁴⁹, M.F. Watson¹⁷, G. Watts¹³⁷, S. Watts⁸¹, A.T. Waugh¹⁴⁹, B.M. Waugh⁷⁶, M. Weber¹²⁸, M.S. Weber¹⁶, P. Weber⁵³, A.R. Weidberg¹¹⁷, P. Weigell⁹⁸, J. Weingarten⁵³, C. Weiser⁴⁷, H. Wellenstein²², P.S. Wells²⁹, T. Wenaus²⁴, D. Wendland¹⁵, Z. Weng^{150,w}, T. Wengler²⁹, S. Wenig²⁹, N. Wermes²⁰, M. Werner⁴⁷, P. Werner²⁹, M. Werth¹⁶², M. Wessels^{57a}, J. Wetter¹⁶⁰, C. Weydert⁵⁴, K. Whalen²⁸, S.J. Wheeler-Ellis¹⁶², A. White⁷, M.J. White⁸⁵, S. White^{121a,121b}, S.R. Whitehead¹¹⁷, D. Whiteson¹⁶², D. Whittington⁵⁹, F. Wicek¹¹⁴, D. Wicke¹⁷⁴, F.J. Wickens¹²⁸, W. Wiedenmann¹⁷², M. Wielers¹²⁸, P. Wienemann²⁰, C. Wiglesworth⁷⁴, L.A.M. Wiik-Fuchs⁴⁷, P.A. Wijeratne⁷⁶, A. Wildauer¹⁶⁶, M.A. Wildt^{41,s}, I. Wilhelm¹²⁵, H.G. Wilkens²⁹, J.Z. Will⁹⁷, E. Williams³⁴, H.H. Williams¹¹⁹, W. Willis³⁴, S. Willocq⁸³, J.A. Wilson¹⁷, M.G. Wilson¹⁴², A. Wilson⁸⁶, I. Wingerter-Seez⁴, S. Winkelmann⁴⁷, F. Winklmeier²⁹, M. Wittgen¹⁴², S.J. Wollstadt⁸⁰, M.W. Wolter³⁸, H. Wolters^{123a,h}, W.C. Wong⁴⁰, G. Wooden⁸⁶, B.K. Wosiek³⁸, J. Wotschack²⁹, M.J. Woudstra⁸¹, K.W. Wozniak³⁸, K. Wraight⁵², C. Wright⁵², M. Wright⁵², B. Wrona⁷², S.L. Wu¹⁷², X. Wu⁴⁸, Y. Wu^{32b,al}, E. Wulf³⁴, B.M. Wynne⁴⁵, S. Xella³⁵, M. Xiao¹³⁵, S. Xie⁴⁷, C. Xu^{32b,z}, D. Xu¹³⁸, B. Yabsley¹⁴⁹, S. Yacoob^{144b}, M. Yamada⁶⁴, H. Yamaguchi¹⁵⁴, A. Yamamoto⁶⁴, K. Yamamoto⁶², S. Yamamoto¹⁵⁴, T. Yamamura¹⁵⁴, T. Yamanaka¹⁵⁴, J. Yamaoka⁴⁴, T. Yamazaki¹⁵⁴, Y. Yamazaki⁶⁵, Z. Yan²¹, H. Yang⁸⁶, U.K. Yang⁸¹, Y. Yang⁵⁹, Z. Yang^{145a,145b}, S. Yanush⁹⁰, L. Yao^{32a}, Y. Yao¹⁴, Y. Yasu⁶⁴, G.V. Ybeles Smit¹²⁹, J. Ye³⁹, S. Ye²⁴, M. Yilmaz^{3c}, R. Yoosoofmiya¹²², K. Yorita¹⁷⁰, R. Yoshida⁵, C. Young¹⁴², C.J. Young¹¹⁷, S. Youssef²¹, D. Yu²⁴, J. Yu⁷, J. Yu¹¹¹, L. Yuan⁶⁵, A. Yurkewicz¹⁰⁵, M. Byszewski²⁹, B. Zabinski³⁸, R. Zaidan⁶¹, A.M. Zaitsev¹²⁷, Z. Zajacova²⁹, L. Zanello^{131a,131b}, A. Zaytsev¹⁰⁶, C. Zeitnitz¹⁷⁴, M. Zeman¹²⁴, A. Zemla³⁸, C. Zender²⁰, O. Zenin¹²⁷, T. Ženiš^{143a}, Z. Zinonos^{121a,121b}, S. Zenz¹⁴, D. Zerwas¹¹⁴, G. Zevi della Porta⁵⁶, Z. Zhan^{32d}, D. Zhang^{32b,ak}, H. Zhang⁸⁷, J. Zhang⁵, X. Zhang^{32d}, Z. Zhang¹¹⁴, L. Zhao¹⁰⁷, T. Zhao¹³⁷, Z. Zhao^{32b}, A. Zhemchugov⁶³, J. Zhong¹¹⁷, B. Zhou⁸⁶, N. Zhou¹⁶², Y. Zhou¹⁵⁰, C.G. Zhu^{32d}, H. Zhu⁴¹, J. Zhu⁸⁶, Y. Zhu^{32b}, X. Zhuang⁹⁷, V. Zhuravlov⁹⁸, D. Zieminska⁵⁹, N.I. Zimin⁶³, R. Zimmermann²⁰, S. Zimmermann²⁰, S. Zimmermann⁴⁷, M. Ziolkowski¹⁴⁰, R. Zitoun⁴, L. Živković³⁴, V.V. Zmouchko^{127,*}, G. Zobernig¹⁷², A. Zoccoli^{19a,19b}, M. zur Nedden¹⁵, V. Zutshi¹⁰⁵, L. Zwalinski²⁹

¹ Physics Department, SUNY Albany, Albany, NY, United States

² Department of Physics, University of Alberta, Edmonton, AB, Canada

³ (a) Department of Physics, Ankara University, Ankara; (b) Department of Physics, Dumlupinar University, Kutahya; (c) Department of Physics, Gazi University, Ankara; (d) Division of Physics, TOBB University of Economics and Technology, Ankara; (e) Turkish Atomic Energy Authority, Ankara, Turkey

⁴ LAPP, CNRS/IN2P3 and Université de Savoie, Annecy-le-Vieux, France

⁵ High Energy Physics Division, Argonne National Laboratory, Argonne, IL, United States

⁶ Department of Physics, University of Arizona, Tucson, AZ, United States

⁷ Department of Physics, The University of Texas at Arlington, Arlington, TX, United States

⁸ Physics Department, University of Athens, Athens, Greece

⁹ Physics Department, National Technical University of Athens, Zografou, Greece

¹⁰ Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan

¹¹ Institut de Física d'Altes Energies and Departament de Física de la Universitat Autònoma de Barcelona and ICREA, Barcelona, Spain

¹² (a) Institute of Physics, University of Belgrade, Belgrade; (b) Vinca Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia

¹³ Department for Physics and Technology, University of Bergen, Bergen, Norway

¹⁴ Physics Division, Lawrence Berkeley National Laboratory and University of California, Berkeley, CA, United States

¹⁵ Department of Physics, Humboldt University, Berlin, Germany

¹⁶ Albert Einstein Center for Fundamental Physics and Laboratory for High Energy Physics, University of Bern, Bern, Switzerland

¹⁷ School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom

¹⁸ (a) Department of Physics, Bogazici University, Istanbul; (b) Division of Physics, Dogus University, Istanbul; (c) Department of Physics Engineering, Gaziantep University, Gaziantep;

(d) Department of Physics, Istanbul Technical University, Istanbul, Turkey

¹⁹ (a) INFN Sezione di Bologna; (b) Dipartimento di Fisica, Università di Bologna, Bologna, Italy

²⁰ Physikalisches Institut, University of Bonn, Bonn, Germany

²¹ Department of Physics, Boston University, Boston, MA, United States

²² Department of Physics, Brandeis University, Waltham, MA, United States

²³ (a) Universidade Federal do Rio De Janeiro COPPE/EE/IF, Rio de Janeiro; (b) Federal University of Juiz de Fora (UFJF), Juiz de Fora; (c) Federal University of Sao Joao del Rei (UFSJ), Sao Joao del Rei; (d) Instituto de Fisica, Universidade de Sao Paulo, Sao Paulo, Brazil

²⁴ Physics Department, Brookhaven National Laboratory, Upton, NY, United States

- 25 ^(a) National Institute of Physics and Nuclear Engineering, Bucharest; ^(b) University Politehnica Bucharest, Bucharest; ^(c) West University in Timisoara, Timisoara, Romania
- 26 Departamento de Física, Universidad de Buenos Aires, Buenos Aires, Argentina
- 27 Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom
- 28 Department of Physics, Carleton University, Ottawa, ON, Canada
- 29 CERN, Geneva, Switzerland
- 30 Enrico Fermi Institute, University of Chicago, Chicago, IL, United States
- 31 ^(a) Departamento de Física, Pontificia Universidad Católica de Chile, Santiago; ^(b) Departamento de Física, Universidad Técnica Federico Santa María, Valparaíso, Chile
- 32 ^(a) Institute of High Energy Physics, Chinese Academy of Sciences, Beijing; ^(b) Department of Modern Physics, University of Science and Technology of China, Anhui; ^(c) Department of Physics, Nanjing University, Jiangsu; ^(d) School of Physics, Shandong University, Shandong, China
- 33 Laboratoire de Physique Corpusculaire, Clermont Université and Université Blaise Pascal and CNRS/IN2P3, Aubiere Cedex, France
- 34 Nevis Laboratory, Columbia University, Irvington, NY, United States
- 35 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- 36 ^(a) INFN Gruppo Collegato di Cosenza; ^(b) Dipartimento di Fisica, Università della Calabria, Arcavata di Rende, Italy
- 37 AGH University of Science and Technology, Faculty of Physics and Applied Computer Science, Krakow, Poland
- 38 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland
- 39 Physics Department, Southern Methodist University, Dallas, TX, United States
- 40 Physics Department, University of Texas at Dallas, Richardson, TX, United States
- 41 DESY, Hamburg and Zeuthen, Germany
- 42 Institut für Experimentelle Physik IV, Technische Universität Dortmund, Dortmund, Germany
- 43 Institut für Kern- und Teilchenphysik, Technical University Dresden, Dresden, Germany
- 44 Department of Physics, Duke University, Durham, NC, United States
- 45 SUPA – School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom
- 46 INFN Laboratori Nazionali di Frascati, Frascati, Italy
- 47 Fakultät für Mathematik und Physik, Albert-Ludwigs-Universität, Freiburg, Germany
- 48 Section de Physique, Université de Genève, Geneva, Switzerland
- 49 ^(a) INFN Sezione di Genova; ^(b) Dipartimento di Fisica, Università di Genova, Genova, Italy
- 50 ^(a) E. Andronikashvili Institute of Physics, Tbilisi State University, Tbilisi; ^(b) High Energy Physics Institute, Tbilisi State University, Tbilisi, Georgia
- 51 II Physikalisches Institut, Justus-Liebig-Universität Giessen, Giessen, Germany
- 52 SUPA – School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom
- 53 II Physikalisches Institut, Georg-August-Universität, Göttingen, Germany
- 54 Laboratoire de Physique Subatomique et de Cosmologie, Université Joseph Fourier and CNRS/IN2P3 and Institut National Polytechnique de Grenoble, Grenoble, France
- 55 Department of Physics, Hampton University, Hampton, VA, United States
- 56 Laboratory for Particle Physics and Cosmology, Harvard University, Cambridge, MA, United States
- 57 ^(a) Kirchhoff-Institut für Physik, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(b) Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg; ^(c) ZITI Institut für Technische Informatik, Ruprecht-Karls-Universität Heidelberg, Mannheim, Germany
- 58 Faculty of Applied Information Science, Hiroshima Institute of Technology, Hiroshima, Japan
- 59 Department of Physics, Indiana University, Bloomington, IN, United States
- 60 Institut für Astro- und Teilchenphysik, Leopold-Franzens-Universität, Innsbruck, Austria
- 61 University of Iowa, Iowa City, IA, United States
- 62 Department of Physics and Astronomy, Iowa State University, Ames, IA, United States
- 63 Joint Institute for Nuclear Research, JINR Dubna, Dubna, Russia
- 64 KEK, High Energy Accelerator Research Organization, Tsukuba, Japan
- 65 Graduate School of Science, Kobe University, Kobe, Japan
- 66 Faculty of Science, Kyoto University, Kyoto, Japan
- 67 Kyoto University of Education, Kyoto, Japan
- 68 Department of Physics, Kyushu University, Fukuoka, Japan
- 69 Instituto de Física La Plata, Universidad Nacional de La Plata and CONICET, La Plata, Argentina
- 70 Physics Department, Lancaster University, Lancaster, United Kingdom
- 71 ^(a) INFN Sezione di Lecce; ^(b) Dipartimento di Matematica e Fisica, Università del Salento, Lecce, Italy
- 72 Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom
- 73 Department of Physics, Jožef Stefan Institute and University of Ljubljana, Ljubljana, Slovenia
- 74 School of Physics and Astronomy, Queen Mary University of London, London, United Kingdom
- 75 Department of Physics, Royal Holloway University of London, Surrey, United Kingdom
- 76 Department of Physics and Astronomy, University College London, London, United Kingdom
- 77 Laboratoire de Physique Nucléaire et de Hautes Energies, UPMC and Université Paris-Diderot and CNRS/IN2P3, Paris, France
- 78 Fysiska institutionen, Lunds universitet, Lund, Sweden
- 79 Departamento de Física Teórica C-15, Universidad Autónoma de Madrid, Madrid, Spain
- 80 Institut für Physik, Universität Mainz, Mainz, Germany
- 81 School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom
- 82 CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France
- 83 Department of Physics, University of Massachusetts, Amherst, MA, United States
- 84 Department of Physics, McGill University, Montreal, QC, Canada
- 85 School of Physics, University of Melbourne, Victoria, Australia
- 86 Department of Physics, The University of Michigan, Ann Arbor, MI, United States
- 87 Department of Physics and Astronomy, Michigan State University, East Lansing, MI, United States
- 88 ^(a) INFN Sezione di Milano; ^(b) Dipartimento di Fisica, Università di Milano, Milano, Italy
- 89 B.I. Stepanov Institute of Physics, National Academy of Sciences of Belarus, Minsk, Belarus
- 90 National Scientific and Educational Centre for Particle and High Energy Physics, Minsk, Belarus
- 91 Department of Physics, Massachusetts Institute of Technology, Cambridge, MA, United States
- 92 Group of Particle Physics, University of Montreal, Montreal, QC, Canada
- 93 P.N. Lebedev Institute of Physics, Academy of Sciences, Moscow, Russia
- 94 Institute for Theoretical and Experimental Physics (ITEP), Moscow, Russia
- 95 Moscow Engineering and Physics Institute (MEPhI), Moscow, Russia
- 96 Skobel'syn Institute of Nuclear Physics, Lomonosov Moscow State University, Moscow, Russia
- 97 Fakultät für Physik, Ludwig-Maximilians-Universität München, München, Germany
- 98 Max-Planck-Institut für Physik (Werner-Heisenberg-Institut), München, Germany
- 99 Nagasaki Institute of Applied Science, Nagasaki, Japan
- 100 Graduate School of Science and Kobayashi-Maskawa Institute, Nagoya University, Nagoya, Japan
- 101 ^(a) INFN Sezione di Napoli; ^(b) Dipartimento di Scienze Fisiche, Università di Napoli, Napoli, Italy

- ¹⁰² Department of Physics and Astronomy, University of New Mexico, Albuquerque, NM, United States
- ¹⁰³ Institute for Mathematics, Astrophysics and Particle Physics, Radboud University Nijmegen/Nikhef, Nijmegen, Netherlands
- ¹⁰⁴ Nikhef National Institute for Subatomic Physics and University of Amsterdam, Amsterdam, Netherlands
- ¹⁰⁵ Department of Physics, Northern Illinois University, DeKalb, IL, United States
- ¹⁰⁶ Budker Institute of Nuclear Physics, SB RAS, Novosibirsk, Russia
- ¹⁰⁷ Department of Physics, New York University, New York, NY, United States
- ¹⁰⁸ Ohio State University, Columbus, OH, United States
- ¹⁰⁹ Faculty of Science, Okayama University, Okayama, Japan
- ¹¹⁰ Homer L. Dodge Department of Physics and Astronomy, University of Oklahoma, Norman, OK, United States
- ¹¹¹ Department of Physics, Oklahoma State University, Stillwater, OK, United States
- ¹¹² Palacký University, RCPTM, Olomouc, Czech Republic
- ¹¹³ Center for High Energy Physics, University of Oregon, Eugene, OR, United States
- ¹¹⁴ LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France
- ¹¹⁵ Graduate School of Science, Osaka University, Osaka, Japan
- ¹¹⁶ Department of Physics, University of Oslo, Oslo, Norway
- ¹¹⁷ Department of Physics, Oxford University, Oxford, United Kingdom
- ¹¹⁸ ^(a) INFN Sezione di Pavia; ^(b) Dipartimento di Fisica, Università di Pavia, Pavia, Italy
- ¹¹⁹ Department of Physics, University of Pennsylvania, Philadelphia, PA, United States
- ¹²⁰ Petersburg Nuclear Physics Institute, Gatchina, Russia
- ¹²¹ ^(a) INFN Sezione di Pisa; ^(b) Dipartimento di Fisica E. Fermi, Università di Pisa, Pisa, Italy
- ¹²² Department of Physics and Astronomy, University of Pittsburgh, Pittsburgh, PA, United States
- ¹²³ ^(a) Laboratório de Instrumentação e Física Experimental de Partículas – LIP, Lisboa, Portugal; ^(b) Departamento de Física Teórica y del Cosmos and CAFPE, Universidad de Granada, Granada, Spain
- ¹²⁴ Institute of Physics, Academy of Sciences of the Czech Republic, Praha, Czech Republic
- ¹²⁵ Faculty of Mathematics and Physics, Charles University in Prague, Praha, Czech Republic
- ¹²⁶ Czech Technical University in Prague, Praha, Czech Republic
- ¹²⁷ State Research Center Institute for High Energy Physics, Protvino, Russia
- ¹²⁸ Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom
- ¹²⁹ Physics Department, University of Regina, Regina, SK, Canada
- ¹³⁰ Ritsumeikan University, Kusatsu, Shiga, Japan
- ¹³¹ ^(a) INFN Sezione di Roma I; ^(b) Dipartimento di Fisica, Università La Sapienza, Roma, Italy
- ¹³² ^(a) INFN Sezione di Roma Tor Vergata; ^(b) Dipartimento di Fisica, Università di Roma Tor Vergata, Roma, Italy
- ¹³³ ^(a) INFN Sezione di Roma Tre; ^(b) Dipartimento di Fisica, Università Roma Tre, Roma, Italy
- ¹³⁴ ^(a) Faculté des Sciences Ain Chock, Réseau Universitaire de Physique des Hautes Energies, Université Hassan II, Casablanca; ^(b) Centre National de l'Energie des Sciences Techniques Nucleaires, Rabat; ^(c) Faculté des Sciences Semlalia, Université Cadi Ayyad, LPHEA, Marrakech; ^(d) Faculté des Sciences, Université Mohamed Premier and LPTPM, Oujda; ^(e) Faculté des Sciences, Université Mohammed V, Agdal, Rabat, Morocco
- ¹³⁵ DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat à l'Energie Atomique), Gif-sur-Yvette, France
- ¹³⁶ Santa Cruz Institute for Particle Physics, University of California Santa Cruz, Santa Cruz, CA, United States
- ¹³⁷ Department of Physics, University of Washington, Seattle, WA, United States
- ¹³⁸ Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom
- ¹³⁹ Department of Physics, Shinshu University, Nagano, Japan
- ¹⁴⁰ Fachbereich Physik, Universität Siegen, Siegen, Germany
- ¹⁴¹ Department of Physics, Simon Fraser University, Burnaby, BC, Canada
- ¹⁴² SLAC National Accelerator Laboratory, Stanford, CA, United States
- ¹⁴³ ^(a) Faculty of Mathematics, Physics & Informatics, Comenius University, Bratislava; ^(b) Department of Subnuclear Physics, Institute of Experimental Physics of the Slovak Academy of Sciences, Kosice, Slovak Republic
- ¹⁴⁴ ^(a) Department of Physics, University of Johannesburg, Johannesburg; ^(b) School of Physics, University of the Witwatersrand, Johannesburg, South Africa
- ¹⁴⁵ ^(a) Department of Physics, Stockholm University; ^(b) The Oskar Klein Centre, Stockholm, Sweden
- ¹⁴⁶ Physics Department, Royal Institute of Technology, Stockholm, Sweden
- ¹⁴⁷ Departments of Physics & Astronomy and Chemistry, Stony Brook University, Stony Brook, NY, United States
- ¹⁴⁸ Department of Physics and Astronomy, University of Sussex, Brighton, United Kingdom
- ¹⁴⁹ School of Physics, University of Sydney, Sydney, Australia
- ¹⁵⁰ Institute of Physics, Academia Sinica, Taipei, Taiwan
- ¹⁵¹ Department of Physics, Technion: Israel Institute of Technology, Haifa, Israel
- ¹⁵² Raymond and Beverly Sackler School of Physics and Astronomy, Tel Aviv University, Tel Aviv, Israel
- ¹⁵³ Department of Physics, Aristotle University of Thessaloniki, Thessaloniki, Greece
- ¹⁵⁴ International Center for Elementary Particle Physics and Department of Physics, The University of Tokyo, Tokyo, Japan
- ¹⁵⁵ Graduate School of Science and Technology, Tokyo Metropolitan University, Tokyo, Japan
- ¹⁵⁶ Department of Physics, Tokyo Institute of Technology, Tokyo, Japan
- ¹⁵⁷ Department of Physics, University of Toronto, Toronto, ON, Canada
- ¹⁵⁸ ^(a) TRIUMF, Vancouver, BC; ^(b) Department of Physics and Astronomy, York University, Toronto, ON, Canada
- ¹⁵⁹ Institute of Pure and Applied Sciences, University of Tsukuba, 1-1-1 Tennodai, Tsukuba, Ibaraki 305-8571, Japan
- ¹⁶⁰ Science and Technology Center, Tufts University, Medford, MA, United States
- ¹⁶¹ Centro de Investigaciones, Universidad Antonio Narino, Bogota, Colombia
- ¹⁶² Department of Physics and Astronomy, University of California Irvine, Irvine, CA, United States
- ¹⁶³ ^(a) INFN Gruppo Collegato di Udine; ^(b) ICTP, Trieste; ^(c) Dipartimento di Chimica, Fisica e Ambiente, Università di Udine, Udine, Italy
- ¹⁶⁴ Department of Physics, University of Illinois, Urbana, IL, United States
- ¹⁶⁵ Department of Physics and Astronomy, University of Uppsala, Uppsala, Sweden
- ¹⁶⁶ Instituto de Física Corpuscular (IFIC) and Departamento de Física Atómica, Molecular y Nuclear and Departamento de Ingeniería Electrónica and Instituto de Microelectrónica de Barcelona (IMB-CNM), University of Valencia and CSIC, Valencia, Spain
- ¹⁶⁷ Department of Physics, University of British Columbia, Vancouver, BC, Canada
- ¹⁶⁸ Department of Physics and Astronomy, University of Victoria, Victoria, BC, Canada
- ¹⁶⁹ Department of Physics, University of Warwick, Coventry, United Kingdom
- ¹⁷⁰ Waseda University, Tokyo, Japan
- ¹⁷¹ Department of Particle Physics, The Weizmann Institute of Science, Rehovot, Israel
- ¹⁷² Department of Physics, University of Wisconsin, Madison, WI, United States
- ¹⁷³ Fakultät für Physik und Astronomie, Julius-Maximilians-Universität, Würzburg, Germany
- ¹⁷⁴ Fachbereich C Physik, Bergische Universität Wuppertal, Wuppertal, Germany
- ¹⁷⁵ Department of Physics, Yale University, New Haven, CT, United States

¹⁷⁶ Yerevan Physics Institute, Yerevan, Armenia

¹⁷⁷ Domaine Scientifique de la Doua, Centre de Calcul CNRS/IN2P3, Villeurbanne Cedex, France

^a Also at Laboratorio de Instrumentacao e Fisica Experimental de Particulas – LIP, Lisboa, Portugal.

^b Also at Faculdade de Ciencias and CFNUL, Universidade de Lisboa, Lisboa, Portugal.

^c Also at Particle Physics Department, Rutherford Appleton Laboratory, Didcot, United Kingdom.

^d Also at TRIUMF, Vancouver, BC, Canada.

^e Also at Department of Physics, California State University, Fresno, CA, United States.

^f Also at Novosibirsk State University, Novosibirsk, Russia.

^g Also at Fermilab, Batavia, IL, United States.

^h Also at Department of Physics, University of Coimbra, Coimbra, Portugal.

ⁱ Also at Department of Physics, UASLP, San Luis Potosi, Mexico.

^j Also at Università di Napoli Parthenope, Napoli, Italy.

^k Also at Institute of Particle Physics (IPP), Canada.

^l Also at Department of Physics, Middle East Technical University, Ankara, Turkey.

^m Also at Louisiana Tech University, Ruston, LA, United States.

ⁿ Also at Dep Fisica and CEFITEC of Faculdade de Ciencias e Tecnologia, Universidade Nova de Lisboa, Caparica, Portugal.

^o Also at Department of Physics and Astronomy, University College London, London, United Kingdom.

^p Also at Group of Particle Physics, University of Montreal, Montreal, QC, Canada.

^q Also at Department of Physics, University of Cape Town, Cape Town, South Africa.

^r Also at Institute of Physics, Azerbaijan Academy of Sciences, Baku, Azerbaijan.

^s Also at Institut für Experimentalphysik, Universität Hamburg, Hamburg, Germany.

^t Also at Manhattan College, New York, NY, United States.

^u Also at School of Physics, Shandong University, Shandong, China.

^v Also at CPPM, Aix-Marseille Université and CNRS/IN2P3, Marseille, France.

^w Also at School of Physics and Engineering, Sun Yat-sen University, Guanzhou, China.

^x Also at Academia Sinica Grid Computing, Institute of Physics, Academia Sinica, Taipei, Taiwan.

^y Also at Dipartimento di Fisica, Università La Sapienza, Roma, Italy.

^z Also at DSM/IRFU (Institut de Recherches sur les Lois Fondamentales de l'Univers), CEA Saclay (Commissariat à l'Energie Atomique), Gif-sur-Yvette, France.

^{aa} Also at Section de Physique, Université de Genève, Geneva, Switzerland.

^{ab} Also at Departamento de Fisica, Universidade de Minho, Braga, Portugal.

^{ac} Also at Department of Physics and Astronomy, University of South Carolina, Columbia, SC, United States.

^{ad} Also at Institute for Particle and Nuclear Physics, Wigner Research Centre for Physics, Budapest, Hungary.

^{ae} Also at California Institute of Technology, Pasadena, CA, United States.

^{af} Also at Institute of Physics, Jagiellonian University, Krakow, Poland.

^{ag} Also at LAL, Université Paris-Sud and CNRS/IN2P3, Orsay, France.

^{ah} Also at Nevis Laboratory, Columbia University, Irvington, NY, United States.

^{ai} Also at Department of Physics and Astronomy, University of Sheffield, Sheffield, United Kingdom.

^{aj} Also at Department of Physics, Oxford University, Oxford, United Kingdom.

^{ak} Also at Institute of Physics, Academia Sinica, Taipei, Taiwan.

^{al} Also at Department of Physics, The University of Michigan, Ann Arbor, MI, United States.

* Deceased.