

Kinetic Theory of Simple Reacting Spheres I

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Abstract

We consider physical and mathematical aspects of the model of *simple reacting spheres* (SRS) in the kinetic theory of chemically reacting fluids. The SRS, being a natural extension of the hard-sphere collisional model, reduces itself to the revised Enskog theory when the chemical reactions are turned off. In the dilute-gas limit, it provides an interesting kinetic model of chemical reactions that has not been considered before. In contrast to other reactive kinetic theories (e.g., *line-of-centers* models), the SRS has built-in *detailed balance* and *microscopic reversibility* conditions. The mathematical analysis of the work consists of global existence result for the system of partial differential equations for the model of SRS.

1 Simple Reacting Spheres

Simple reacting spheres (SRS) has been developed by N. Xystris, J. S. Dahler [1] and further advanced by J. S. Dahler and L. Quin in [2], [3]. The present paper is the first in a series of our articles on physical and mathematical properties of SRS. In the SRS model, the molecules behave as if they were single mass points with two internal states of excitation. Collisions may alter the internal states: this occurs when the kinetic energy associated with the reactive motion exceeds the activation energy. Reactive and non-reactive collision events are considered to be hard spheres-like. In a four component mixture A, B, A^*, B^* , the chemical reactions are of the type: $A + B \rightleftharpoons A^* + B^*$. Here, A^* and B^* are distinct species from A and B . We use the indices 1, 2, 3, and 4 for the particles A, B, A^* , and B^* respectively. Furthermore, m_i and d_i denote the mass and the diameter of the i -th particle, $i = 1, \dots, 4$, and reactions take place when the reactive particles are separated by a distance $\sigma_{12} = \frac{1}{2}(d_1 + d_2)$ or $\sigma_{34} = \frac{1}{2}(d_3 + d_4)$. The conservation of mass has the form $m_1 + m_2 = m_3 + m_4 = M$. Reactions take place when the reactive particles are separated by a distance $\sigma_{12} = \frac{1}{2}(d_1 + d_2)$, where d_i denotes the diameter of the i -th particle.

1.1 Elastic encounters

In the case of elastic collisions between a pair of particles from species i and s , the initial velocities v, w take post-collisional values

$$v' = v - 2\frac{\mu_{is}}{m_i}\epsilon\langle\epsilon, v - w\rangle, \quad w' = w + 2\frac{\mu_{is}}{m_s}\epsilon\langle\epsilon, v - w\rangle. \quad (1)$$

Here, $\langle \cdot, \cdot \rangle$ is the inner product in \mathbb{R}^3 , ϵ is a vector along the line passing through the centers of the spheres at the moment of impact, i.e., $\epsilon \in \mathbb{S}_+^2 = \{\epsilon \in \mathbb{R}^3 : |\epsilon| = 1, \langle \epsilon, v - w \rangle \geq 0\}$

and $\mu_{is} = m_i m_s / (m_i + m_s)$ is the reduced mass of the colliding pair. m_i and m_s are the masses of particles from i -th and s -th species, respectively ($i, s = 1, 2, 3, 4$).

1.2 Reactive encounters

For the reactive collision between particles from species i and s to occur ($i, s = 1, \dots, 4$), the kinetic energy associated with the relative motion along the line of centers must exceed the activation energy γ_i ,

$$(1/2)\mu_{is}(\langle \epsilon, v - w \rangle)^2 \geq \gamma_i, \quad (2)$$

In the case of the (endothermic) reaction $A + B \rightarrow A^* + B^*$ the velocities v, w take their post-reactive values

$$v^\dagger = \frac{1}{M} \left[m_1 v + m_2 w + m_4 \sqrt{\frac{\mu_{12}}{\mu_{34}}} \left\{ (v - w) - \epsilon \langle \epsilon, v - w \rangle + \epsilon \alpha^- \right\} \right], \quad (3)$$

$$w^\dagger = \frac{1}{M} \left[m_1 v + m_2 w - m_3 \sqrt{\frac{\mu_{12}}{\mu_{34}}} \left\{ (v - w) - \epsilon \langle \epsilon, v - w \rangle + \epsilon \alpha^- \right\} \right], \quad (4)$$

with $\alpha^- = \sqrt{(\langle \epsilon, v - w \rangle)^2 - 2E_{abs}/\mu_{12}}$, and E_{abs} the energy absorbed by the internal degrees of freedom. The absorbed energy E_{abs} has the property $E_{abs} = E_3 + E_4 - E_1 - E_2 > 0$, where $E_i > 0$, $i = 1, \dots, 4$, is the energy of i -th particle associated with its internal degrees of freedom.

Now, in order to complete the definition of the model, the activation energies γ_1, γ_2 for A and B are chosen to satisfy $\gamma_1 \geq E_{abs} > 0$, and by symmetry, $\gamma_2 = \gamma_1$.

For the inverse (exothermic) reaction, $A^* + B^* \rightarrow A + B$, the post-reactive velocities are given by

$$v^\dagger = \frac{1}{M} \left[m_3 v + m_4 w + m_2 \sqrt{\frac{\mu_{34}}{\mu_{12}}} \left\{ (v - w) - \epsilon \langle \epsilon, v - w \rangle + \epsilon \alpha^+ \right\} \right], \quad (5)$$

$$w^\dagger = \frac{1}{M} \left[m_3 v + m_4 w - m_1 \sqrt{\frac{\mu_{34}}{\mu_{12}}} \left\{ (v - w) - \epsilon \langle \epsilon, v - w \rangle + \epsilon \alpha^+ \right\} \right], \quad (6)$$

with $\alpha^+ = \sqrt{(\langle \epsilon, v - w \rangle)^2 + 2E_{abs}/\mu_{34}}$, and the activation energies for A^* and B^* being $\gamma_3 = \gamma_1 - E_{abs}$ and $\gamma_4 = \gamma_3$.

Post- and pre-collisional velocities of the reactive pairs satisfy conservation of the momentum

$$m_1 v + m_2 w = m_3 v^\dagger + m_4 w^\dagger, \quad m_3 v + m_4 w = m_1 v^\dagger + m_2 w^\dagger. \quad (7)$$

A part of kinetic energy is exchanged with the energy absorbed by the internal states. The following equalities hold:

$$m_1 v^2 + m_2 w^2 = m_3 v^{\dagger 2} + m_4 w^{\dagger 2} + 2E_{abs}, \quad m_3 v^2 + m_4 w^2 = m_1 v^{\dagger 2} + m_2 w^{\dagger 2} - 2E_{abs}. \quad (8)$$

1.3 The system of equations

For $i = 1, 2, 3, 4$, $f_i(t, x, v)$ denotes the one-particle distribution function of the i th component of the reactive mixture. The function $f_i(t, x, v)$, which changes in time due to free

streaming and collisions (elastic and reactive), represents at time t the number density of particles at point x with velocity v .

The SRS kinetic system has the form

$$\frac{\partial f_i}{\partial t} + v \frac{\partial f_i}{\partial x} = J_i^E + J_i^R, \quad i = 1, 2, 3, 4, \quad (9)$$

where J_i^E is the non-reactive (hard-sphere) collision operator

$$\begin{aligned} J_i^E = & \sum_{s=1}^4 \left\{ \sigma_{is}^2 \iint_{\mathbb{R}^3 \times \mathbb{S}_+^2} \left[f_{is}^{(2)}(t, x, v', x - \sigma_{is}\epsilon, w') - f_{is}^{(2)}(t, x, v, x + \sigma_{is}\epsilon, w) \right] \langle \epsilon, v - w \rangle d\epsilon dw \right\} \\ & - \beta_{ij} \sigma_{ij}^2 \iint_{\mathbb{R}^3 \times \mathbb{S}_+^2} \left[f_{ij}^{(2)}(t, x, v', x - \sigma_{ij}\epsilon, w') - f_{ij}^{(2)}(t, x, v, x + \sigma_{ij}\epsilon, w) \right] \\ & \times \Theta(\langle \epsilon, v - w \rangle - \Gamma_{ij}) \langle \epsilon, v - w \rangle d\epsilon dw, \end{aligned} \quad (10)$$

and $f_{is}^{(2)}$ approximates the density of pairs of particles in collisional configurations. The second term in (10), with β_{ij} in front of it, singles out those pre-collisional states that are energetic enough to result in the reaction, and thus preventing double counting of the events in the collisional integrals. In the case when $\beta_{ij} = 0$, for $i, j = 1, \dots, 4$, the term J_i^E , in (10), reduces to two-particle collisional operator for 4-species mixtures with hard-sphere potential. For $i = 1, 2, 3, 4$, the reactive terms are

$$\begin{aligned} J_i^R = & \beta_{ij} \sigma_{ij}^2 \iint_{\mathbb{R}^3 \times \mathbb{S}_+^2} \left[\left(\frac{\mu_{ij}}{\mu_{kl}} \right)^{3/2} f_{kl}^{(2)}(t, x, v_{ij}^\circ, x - \sigma_{ij}\epsilon, w_{ij}^\circ) - f_{ij}^{(2)}(t, x, v, x + \sigma_{ij}\epsilon, w) \right] \\ & \times \Theta(\langle \epsilon, v - w \rangle - \Gamma_{ij}) \langle \epsilon, v - w \rangle d\epsilon dw. \end{aligned} \quad (11)$$

Here, $0 \leq \beta_{ij} \leq 1$ are the steric factors, $\Gamma_{ij} = \sqrt{2\gamma_i/\mu_{ij}}$, and Θ is the Heaviside step function. The pairs of post-reactive velocities are $(v_{ij}^\circ, w_{ij}^\circ) = (v^\ddagger, w^\ddagger)$ for $i, j = 1, 2$, and $(v_{ij}^\circ, w_{ij}^\circ) = (v^\dagger, w^\dagger)$ for $i, j = 3, 4$. Pairs of indices (i, j) and (k, l) are from the set of quadruples (i, j, k, l) : $\{(1, 2, 3, 4), (2, 1, 4, 3), (3, 4, 1, 2), (4, 3, 2, 1)\}$.

Lemma 1.

(1) For $i, s = 1, 2, 3, 4$, the inverse velocities to v', w' are given by

$$v = v' - 2 \frac{\mu_{is}}{m_i} \epsilon \langle \epsilon, v' - w' \rangle, \quad w = w' + 2 \frac{\mu_{is}}{m_s} \epsilon \langle \epsilon, v' - w' \rangle. \quad (12)$$

For fixed ϵ , the Jacobian of the transformation $(v, w) \mapsto (v', w')$ is equal to -1 . Furthermore, $\langle \epsilon, v' - w' \rangle = -\langle \epsilon, v - w \rangle$,

(2) The inverse velocities to v^\ddagger, w^\ddagger are given by

$$v = \frac{1}{M} \left[m_3 v^\ddagger + m_4 w^\ddagger + m_2 \sqrt{\frac{\mu_{34}}{\mu_{12}}} \left\{ (v^\ddagger - w^\ddagger) - \epsilon \langle \epsilon, v^\ddagger - w^\ddagger \rangle + \epsilon \alpha^+ \right\} \right], \quad (13)$$

$$w = \frac{1}{M} \left[m_3 v^\ddagger + m_4 w^\ddagger - m_1 \sqrt{\frac{\mu_{34}}{\mu_{12}}} \left\{ (v^\ddagger - w^\ddagger) - \epsilon \langle \epsilon, v^\ddagger - w^\ddagger \rangle + \epsilon \alpha^+ \right\} \right], \quad (14)$$

and the inverse velocities to v^\dagger, w^\dagger are given by

$$v = \frac{1}{M} \left[m_1 v^\dagger + m_2 w^\dagger + m_4 \sqrt{\frac{\mu_{12}}{\mu_{34}}} \left\{ (v^\dagger - w^\dagger) - \epsilon \langle \epsilon, v^\dagger - w^\dagger \rangle + \epsilon \alpha^- \right\} \right], \quad (15)$$

$$w = \frac{1}{M} \left[m_1 v^\dagger + m_2 w^\dagger - m_3 \sqrt{\frac{\mu_{12}}{\mu_{34}}} \left\{ (v^\dagger - w^\dagger) - \epsilon \langle \epsilon, v^\dagger - w^\dagger \rangle + \epsilon \alpha^- \right\} \right]. \quad (16)$$

(3) For fixed ϵ , the Jacobians of the transformations $(v, w) \mapsto (v^\dagger, w^\dagger)$ and $(v, w) \mapsto (v^\ddagger, w^\ddagger)$ are given by

$$\left(\frac{\mu_{34}}{\mu_{12}} \right)^{3/2} \frac{\langle \epsilon, v - w \rangle}{\alpha^+} \quad \text{and} \quad \left(\frac{\mu_{12}}{\mu_{34}} \right)^{3/2} \frac{\langle \epsilon, v - w \rangle}{\alpha^-}, \quad (17)$$

respectively.

(4) Furthermore, $\langle \epsilon, v^\dagger - w^\dagger \rangle = \alpha^+$, $\langle \epsilon, v^\ddagger - w^\ddagger \rangle = \alpha^-$, $\frac{1}{2} \mu_{12} (\langle \epsilon, v - w \rangle)^2 - \gamma_1 = \frac{1}{2} \mu_{34} (\langle \epsilon, v^\ddagger - w^\ddagger \rangle)^2 - \gamma_3$, and $\frac{1}{2} \mu_{34} (\langle \epsilon, v - w \rangle)^2 - \gamma_3 = \frac{1}{2} \mu_{12} (\langle \epsilon, v^\dagger - w^\dagger \rangle)^2 - \gamma_1$.

Due to space limitations, the proof of Lemma 1 will be provided in the forthcoming work.

2 The dilute SRS kinetic system

The system of equations (9)-(11) requires a closure relation for $f_{is}^{(2)}$. In the case of moderately dense gases, the two-particle distribution function $f_{is}^{(2)}$ is usually approximated by

$$f_{is}^{(2)}(t, x_1, v_1, x_2, v_2) = g_{is}^{(2)}(x_1, x_2 | \{n_i(t, \cdot)\}) f_i(t, x_1, v_1) f_s(t, x_2, v_2), \quad (18)$$

where $n_i(t, x) = \int_{\mathbb{R}^3} f_i(t, x, v) dv$ is the local number density of the component i and $g_{ij}^{(2)}$ is the known pair correlation function for a non-uniform hard-sphere system at equilibrium with the local densities $n_i(t, x)$. The notation $g_{ij}^{(2)}(x_1, x_2 | \{n_i(t, \cdot)\})$ indicates that $g_{ij}^{(2)}$ is a functional of the local densities n_i . The closure relation (18) is employed in [2] and [3]. Finally, in the case of non-reactive mixtures ($\beta_{ij} = 0$, for $i, j = 1, \dots, 4$), the corresponding system of equations (9)-(11) becomes the revised Enskog system for 4-species mixtures [4]. In this work, we will consider a dilute gas regime with the corresponding closure relation given by:

$$f_{is}^{(2)}(t, x_1, v_1, x_2, v_2) = f_i(t, x_1, v_1) f_s(t, x_2, v_2). \quad (19)$$

The system of equations (9)-(11) takes the form:

$$\frac{\partial f_i}{\partial t} + v \frac{\partial f_i}{\partial x} = J_i^E + J_i^R, \quad f_i(0, x, v) = f_{i0}(x, v), \quad i = 1, \dots, 4, \quad (x, v) \in \Omega \times \mathbb{R}^3, \quad (20)$$

with

$$\begin{aligned} J_i^E &= \sum_{s=1}^4 \left\{ \sigma_{is}^2 \iint_{\mathbb{R}^3 \times \mathbb{S}_+^2} \left[f_i(t, x, v') f_s(t, x, w') - f_i(t, x, v) f_s(t, x, w) \right] \langle \epsilon, v - w \rangle d\epsilon dw \right\} \\ &\quad - \beta_{ij} \sigma_{ij}^2 \iint_{\mathbb{R}^3 \times \mathbb{S}_+^2} \left[f_i(t, x, v') f_j(t, x, w') - f_i(t, x, v) f_j(t, x, w) \right] \\ &\quad \times \Theta(\langle \epsilon, v - w \rangle - \Gamma_{ij}) \langle \epsilon, v - w \rangle d\epsilon dw, \end{aligned} \quad (21)$$

and

$$J_i^R = \beta_{ij} \sigma_{ij}^2 \iint_{\mathbb{R}^3 \times \mathbb{S}_+^2} \left[\left(\frac{\mu_{ij}}{\mu_{kl}} \right)^{3/2} f_k(t, x, v_{ij}^\circ) f_l(t, x, w_{ij}^\circ) - f_i(t, x, v) f_j(t, x, w) \right] \times \Theta(\langle \epsilon, v - w \rangle - \Gamma_{ij}) \langle \epsilon, v - w \rangle d\epsilon dw, \quad (22)$$

where f_{i0} , $i = 1, \dots, 4$ are suitable nonnegative initial conditions that will be defined later and $\Omega \subseteq \mathbb{R}^3$ denotes the spatial domain of the gas mixture. We consider two choices for the set Ω : $\Omega = \mathbb{R}^3$, or Ω being a 3-dimensional torus $[0, L]^3$, $L > 0$. The latter choice corresponds to case of the periodic boundary conditions on $[0, L]^3$. Also, $\Gamma_{ij} = \sqrt{2\gamma_i/\mu_{ij}}$ and Θ is the Heaviside step function. As before, the pairs of post-reactive velocities are $(v_{ij}^\circ, w_{ij}^\circ) = (v^\ddagger, w^\ddagger)$ for $i, j = 1, 2$, and $(v_{ij}^\circ, w_{ij}^\circ) = (v^\dagger, w^\dagger)$ for $i, j = 3, 4$. The pairs of indices (i, j) and (k, l) are from the set of quadruples (i, j, k, l) : $\{(1, 2, 3, 4), (2, 1, 4, 3), (3, 4, 1, 2), (4, 3, 2, 1)\}$.

Proposition 1. *Assume that $\beta_{ij} = \beta_{ji}$ for $(i, j) \in \{(1, 2), (2, 1), (3, 4), (4, 3)\}$. For ϕ_i measurable on $\Omega \times \mathbb{R}^3$ and $f_i \in C_0(\Omega \times \mathbb{R}^3)$, $i = 1, \dots, 4$, we have:*

$$\sum_{i=1}^4 \int_{\mathbb{R}^3} \phi_i J_i^E dv = \sum_{i=1}^4 \sum_{s=1}^4 \sigma_{is}^2 \iiint_{\mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{S}_+^2} [\phi_i(x, v) + \phi_s(x, w) - \phi_i(x, v') - \phi_s(x, w')] \times [f_i(v') f_s(w') - f_i(v) f_s(w)] \langle \epsilon, v - w \rangle \Xi_{is} d\epsilon dw dv, \quad (23)$$

$$\sum_{i=1}^4 \int_{\mathbb{R}^3} \phi_i J_i^R dv = \iiint_{\mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{S}_+^2} [\beta_{12} \sigma_{12}^2 \phi_1(x, v) + \beta_{21} \sigma_{21}^2 \phi_2(x, w) - \beta_{34} \sigma_{34}^2 \phi_3(x, v^\ddagger) - \beta_{43} \sigma_{43}^2 \phi_4(x, w^\ddagger)] \times \left[\left(\frac{\mu_{12}}{\mu_{34}} \right)^{3/2} f_3(x, v^\ddagger) f_4(x, w^\ddagger) - f_1(x, v) f_2(x, w) \right] \Theta(\langle \epsilon, v - w \rangle - \Gamma_{12}) \langle \epsilon, v - w \rangle d\epsilon dw dv, \quad (24)$$

where Ξ_{is} , appearing in (23), is given by

$$\Xi_{is} = \begin{cases} \frac{1}{2} \Theta(\langle \epsilon, v - w \rangle - \Gamma_{is}) + \frac{1}{2} (1 - \beta_{is}) \Theta(\Gamma_{is} - \langle \epsilon, v - w \rangle), & \text{if } (i, s) \in I, \\ \frac{1}{4} \Theta(\langle \epsilon, v - w \rangle), & \text{if } i = s, \\ \frac{1}{2} \Theta(\langle \epsilon, v - w \rangle), & \text{otherwise,} \end{cases} \quad (25)$$

with $I = \{(1, 2), (2, 1), (3, 4), (4, 3)\}$.

The post-collisional velocities, v' and w' , are given in (1), while the post-reactive velocities, v^\ddagger and w^\ddagger , are given in (3)-(4).

Due to space limitations, the proof of Proposition 1 will be provided in the forthcoming work.

3 Conservation laws

Under the additional assumption $\beta_{12}\sigma_{12}^2 = \beta_{34}\sigma_{34}^2$, Proposition 1 implies that for any $a, c \in \mathbb{R}$ and $b \in \mathbb{R}^3$,

$$\phi_i(x, v) = am_i + m_i\langle b, v \rangle + c \left(\frac{m_i v^2}{2} + E_i \right), \quad i = 1, \dots, 4, \quad \implies \begin{cases} \sum_{i=1}^4 \int_{\mathbb{R}^3} \phi_i J_i^E dv = 0, \\ \sum_{i=1}^4 \int_{\mathbb{R}^3} \phi_i J_i^R dv = 0. \end{cases} \quad (26)$$

Property (26) implies that if f_i is a nonnegative smooth solution of (20) on $[0, T]$, $T > 0$, then, at least formally, we have the following conservation laws for $t \in [0, T]$, relative to mass, momentum and total energy:

$$\sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} m_i f_i(t, x, v) dv dx = \sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} m_i f_{i0}(x, v) dv dx, \quad (27)$$

$$\sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} m_i v f_i(t, x, v) dv dx = \sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} m_i v f_{i0}(x, v) dv dx, \quad (28)$$

$$\sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} \left(\frac{m_i v^2}{2} + E_i \right) f_i(t, x, v) dv dx = \sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} \left(\frac{m_i v^2}{2} + E_i \right) f_{i0}(x, v) dv dx, \quad (29)$$

where $f_{i0}(x, v)$, $i = 1, \dots, 4$, are nonnegative initial conditions of the dilute SRS kinetic system (20). For f_i , a smooth solution with compact support, the above conservation laws follow easily from multiplying i -th equation of the dilute SRS system by ϕ_i , integrating with respect to $(t, x, v) \in [0, T] \times \Omega \times \mathbb{R}^3$, and then applying (26).

4 Entropy identity, H -function, and equilibrium solutions

Proposition 1 also implies existence of a Liapunov functional (an H -function) for (20)-(22) consistent with system's physical equilibrium. Assume that for $i, j = 1, \dots, 4$, the conditions $\beta_{ij} = \beta_{ji}$ and $\beta_{12}\sigma_{12}^2 = \beta_{34}\sigma_{34}^2$ are satisfied. For f_i , a smooth nonnegative solution, we multiply (20) by $1 + \log(f_i/\mu_{ij})$ with $i = 1, \dots, 4$ and $(i, j) \in \{(1, 2), (2, 1), (3, 4), (4, 3)\}$, integrate over $\Omega \times \mathbb{R}^3$, and use (23)–(24) (with $\phi_i = \log(f_i/\mu_{ij})$) to obtain the following entropy identity:

$$\begin{aligned} & \frac{d}{dt} \sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} f_i \log(f_i/\mu_{ij}) dv dx + \sum_{i,s=1}^4 \sigma_{is}^2 \int \cdots \int_{\Omega \times \mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{S}_+^2} \left[f_i(v') f_s(w') - f_i(v) f_s(w) \right] \\ & \times \log \left(\frac{f_i(v') f_s(w')}{f_i(v) f_s(w)} \right) \langle \epsilon, v - w \rangle \Xi_{is} dv dw dv dx \end{aligned}$$

$$\begin{aligned}
& + \beta_{12}\sigma_{12}^2 \int \cdots \int_{\Omega \times \mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{S}_+^2} \left\{ \left[\left(\frac{\mu_{12}}{\mu_{34}} \right)^{3/2} f_3(v^\dagger) f_4(w^\dagger) - f_1(v) f_2(w) \right] \times \right. \\
& \left. \log \left[\left(\frac{\mu_{12}}{\mu_{34}} \right)^{3/2} \frac{f_3(v^\dagger) f_4(w^\dagger)}{f_1(v) f_2(w)} \right] \Theta(\langle \epsilon, v - w \rangle - \Gamma_{12}) \langle \epsilon, v - w \rangle \right\} d\epsilon dw dv dx = 0. \quad (30)
\end{aligned}$$

We observe that the second and the third terms in the left hand side of (30) are nonnegative and thus, the convex function $H(t)$, defined by

$$H(t) = \sum_{i=1}^4 \iint_{\Omega \times \mathbb{R}^3} f_i(t, x, v) \log [f_i(t, x, v) / \mu_{ij}] dv dx, \quad (31)$$

is non-increasing in $t \geq 0$. Hence $H(t)$ is an H -function (Liapunov functional) of the system (20)-(22).

For the standard macroscopic quantities (moments of f_i), namely the number densities $n_i(t, x)$, the macroscopic velocity $u(t, x)$, and the macroscopic temperature $\mathcal{T}(t, x)$, we have the following characterization of equilibrium solutions for the system (20)-(22):

Proposition 2. *Assume that for $i, j = 1, \dots, 4$, the coefficients $0 < \beta_{ij} \leq 1$ satisfy the conditions $\beta_{ij} = \beta_{ji}$ and $\beta_{12}\sigma_{12}^2 = \beta_{34}\sigma_{34}^2$. Let $n_i(t, x) \geq 0$, $u(t, x)$, and $\mathcal{T}(t, x) \geq 0$ be given measurable functions. Then for all $0 \leq f_i \in L^1(\Omega \times \mathbb{R}^3)$, the following statements are equivalent:*

1. $f_i = n_i \left(\frac{m_i}{2\pi k\mathcal{T}} \right)^{3/2} \exp \left(-\frac{m_i(v-u)^2}{2k\mathcal{T}} \right)$, $i = 1, \dots, 4$,
and $n_1 n_2 = \left(\frac{\mu_{12}}{\mu_{34}} \right)^{3/2} n_3 n_4 \exp \left(\frac{E_{abs}}{k\mathcal{T}} \right)$;
2. $J_i^E(\{f_i\}) = 0$ and $J_i^R(\{f_i\}) = 0$, $i = 1, \dots, 4$;
3. $\sum_{i=1}^4 \int_{\mathbb{R}^3} [J_i^E(\{f_i\}) + J_i^R(\{f_i\})] \log(f_i / \mu_{ij}) dv = 0$.

The proof of Proposition 2 follows a similar line of arguments as the proof of Proposition 3.2 in [5].

5 Existence result

Definition 1. *A nonnegative $f_i \in L_{loc}^1((0, T) \times \Omega \times \mathbb{R}^3)$, $i = 1, 2, 3, 4$, is a mild solution of the system (20)-(22) if for each $0 < T < \infty$, the gain and loss terms of J_i^E and J_i^R are in $L^1(0, T)$, a.e. (almost everywhere) in $(x, v) \in \Omega \times \mathbb{R}^3$ and*

$$f_i^\#(t, x, v) - f_i^\#(s, x, v) = \int_s^t \left[J_i^E(\{f_i\})^\#(\tau, x, v) + J_i^R(\{f_i\})^\#(\tau, x, v) \right] d\tau, \quad 0 < s < t \leq T, \quad (32)$$

where $f_i^\#(t, x, v) = f(t, x + tv, v)$ and with similar definitions for $J_i^{E\#}$ and $J_i^{R\#}$.

The following result generalizes Theorem 4.2 of [5] (see also Theorem 5.2 of [6]).

Theorem 1. *Assume that for $i, j = 1, \dots, 4$, the coefficients $0 \leq \beta_{ij} \leq 1$ satisfy the conditions $\beta_{ij} = \beta_{ji}$ and $\beta_{12}\sigma_{12}^2 = \beta_{34}\sigma_{34}^2$. If for each $i = 1, \dots, 4$, the initial conditions $f_{i0} \geq 0$, satisfy*

$$\sup_i \iint_{\Omega \times \mathbb{R}^3} (1 + x^2 + v^2 + \log^+ f_{i0}) f_{i0} \, dv dx = C_0 < \infty, \quad (33)$$

with $\log^+(z) = \max\{\log(z), 0\}$, then there exists a nonnegative mild solution $\{f_i\}$ of the system (20)-(22), with $f_i \in C([0, T]; L^1(\Omega \times \mathbb{R}^3))$ and such that $f_i(t)|_{t=0} = f_{i0}$, for $i = 1, 2, 3, 4$.

The proof of Theorem 1 is similar to Theorem 4.2 of [5] and will not be provided here.

Remark 1. *In the case $\beta_{ij} = 0$, for $i, j = 1, \dots, 4$, Theorem 1 provides existence result for chemically inert system (20)-(22).*

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